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## UNSTEADY AERODYNAMICS FOR ADVANCED CONFIGURATIONS

### PART II—A TRANSONIC BOX METHOD FOR PLANAR LIFTING SURFACES

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## FOREWORD

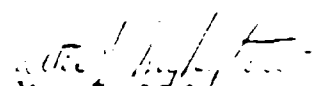
This part of the report covers a portion of the research conducted by the Space and Information Systems Division of North American Aviation, Inc., Downey, California, for the Aerospace Dynamics Branch, Vehicle Dynamics Division, Air Force Flight Dynamics Laboratory, Wright-Patterson Air Force Base, Ohio, under Contract No. AF 33(657)-10399.

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## ABSTRACT

The fundamental equations of the transonic box method were derived, based on the representation of the velocity potential by a doublet distribution. They form the basis of a systematic method of treating an oscillating wing at  $M = 1$ , analogous to the supersonic Mach box method.

A digital computer program, written in Fortran IV, is presented. The program applies to a planar wing of polygonal planform, with a straight trailing edge, and as many as three sweep angles along the leading edge. For a maximum of ten modes of oscillation, the program computes the oscillatory potentials and pressures and a generalized force matrix.

Results obtained from the program are compared with existing theoretical and experimental values. Several possible extensions of the method are described.

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## LIST OF SYMBOLS

<u>Symbol</u>	<u>Definition</u>
$a$	Local speed of sound; speed of sound at infinity
$a_{nm}$	Coefficient in the potential series
$A(i - j,  j - j' )$	Influence coefficient: the upwash at the center of $B_{ij}$ caused by a unit doublet distribution over $B_{i'j'}$
$A_{jr}$	Term in $\bar{\varphi}$ evaluated at $(x_j, y_j)$
$AXY(I, J)$	An integral over the wing planform
$AY(J)$	An integral along the trailing edge
$b$	Root chord length
$B$	Region composed of boxes, approximating $S$
$B_{ij}$	A box
$(B_{rr'})$	Matrix used in least squares surface fits
$BXY$	Part of $AXY(I, J)$
$BY$	Part of $AY(J)$
$\bar{C}_p$	Pressure coefficient
$(C'_r), (C''_r)$	Column matrices used in least squares surface fits
$d$	Dimensionless length of box side
$d_{nm}$	Coefficient in deflection polynomial
$DA$	The data array
$f$	Function which describes the wing deflection
$F$	Factor which gives $\bar{\varphi}$ the proper edge behavior

(Continued on next page.)



<u>Symbol</u>	<u>Definition</u>
$\tilde{g}_u, \tilde{g}_l, \tilde{f}$	Functions used in the equations of upper and lower wing surfaces
$h_j$	Weight used in Gaussian quadrature
$i$	$\sqrt{-1}$
$i, j$	Indexes specifying box position
$I, J$	Indexes
$k$	Reduced frequency: $\omega b / U_\infty$
$l$	$kd$
$L_{ij}$	Generalized force coefficient
$M$	Mach number
$n, m$	Indexes equal to power of $x$ and power of $y^2$
$NC$	Number of coefficients
$NP$	Number of points
$NS$	Number of segments of leading edge given by the data
$p, q$	Integration variables
$Q$	Quantity minimized in least squares surface fits
$r$	Index
$S$	Function used in the equation of a surface
$S$	Region in the $xy$ -plane occupied by the wing; the area of this region
$t$	Time
$u, v$	Integration variables
$u_j$	Point used in Gaussian quadrature
$U_\infty$	Air speed of the wing; speed of flow at infinity

(Continued on next page.)

<u>Symbol</u>	<u>Definition</u>
$w$	Upwash at $z = 0+$
$W$	The region of the $xy$ -plane occupied by the wing's wake
$\tilde{x}, \tilde{y}, \tilde{z}$	Coordinates with dimensions of length
$x, y, z$	Dimensionless coordinates
$(x_i, y_j)$	Center of $B_{ij}$
$(x_j, y_j)$	Point at which a value of potential or deflection is given
$x_1, \dots, x_{NS}$ $y_1, \dots, y_{NS}$	Coordinates of points on the leading edge given by the data
$x_0$	Function which describes the leading edge: $x = x_0(y)$
$y_{max}$	Value of $y$ at the wing tip
$y_+, y_-$	Limits of integration
$a_{nm}, a_r$	Real part of $a_{nm}$
$\beta_r$	Imaginary part of $a_{nm}$
$\delta$	Constant factor in the deflection
$\Delta p_i$	Lifting pressure in the $i$ th mode
$\xi, \eta$	Integration variables equivalent to $x, y$
$\nu$	Frequency
$\rho$	Density
$\rho$	Source or doublet strength
$\sigma$	The integral over $x$ involved in BXY
$\Phi$	Velocity potential
$\phi$	Steady perturbation potential
$\varphi$	Unsteady perturbation potential

(Continued on next page.)

<u>Symbol</u>	<u>Definition</u>
$\bar{\varphi}$	Time independent factor of $\varphi$
$\bar{\varphi}_0$	Potential of a point source
$\bar{\varphi}_1$	Potential of a point doublet
$\bar{\varphi}_s$	Potential of a source distribution
$\bar{\varphi}_d$	Potential of a doublet distribution
$\bar{\varphi}_{ij}$	Value of $\bar{\varphi}$ in $B_{ij}$
$\bar{\varphi}'_j$	Real part of value of $\bar{\varphi}$ at $(x_j, y_j)$
$\Psi$	Upwash in the xy-plane caused by a point doublet
$\omega$	Angular frequency, $2\pi\nu$

## 1. INTRODUCTION

The transonic box program is designed to calculate the unsteady potentials for a given set of modes of wing oscillation and to compute the generalized forces. Pressure distributions may be obtained from the potentials.

A planar wing with a straight trailing edge is assumed. The oscillations are assumed to be symmetric in the spanwise coordinate  $y$ . None of these assumptions is necessary for the method. (See Section 5.)

The basic step in the box method is the solution of the system of simultaneous equations [Equation (33)] which determine a set of values of potential on the wing from a corresponding array of upwash values. A surface is fitted to these values, giving a functional representation of the potential that is used subsequently to find pressures and generalized forces.

The method used is suggested by the success of supersonic box methods (References 1 through 4). The potential is generated by a doublet distribution rather than by a source distribution because the latter method would involve diaphragm regions of infinite extent, whereas the doublet distribution is confined to the wing and its wake.

## 2. THEORETICAL DEVELOPMENT OF THE METHOD

### 1. THE DIFFERENTIAL EQUATION

We consider an oscillating body moving at speed  $U_\infty$  through a nonviscous fluid. From the point of view of a moving coordinate system  $(\bar{x}, \bar{y}, \bar{z})$  in which the average position of the body is fixed, there is a flow past the body with velocity  $U_\infty$  at infinity. Assume that the flow is irrotational; then the velocity field of the flow is the gradient of a potential function  $\Phi$ , which satisfies the differential equation

$$\nabla^2 \Phi - \frac{1}{a^2} \left[ \Phi_{tt} + 2 \nabla \Phi \cdot \nabla \Phi_t + (\nabla \Phi \cdot \nabla) 1/2 (\nabla \Phi)^2 \right] = 0 \quad (1)$$

(See Reference 5, p. 193 where  $a$  is the local speed of sound.

Suppose that the flow is approximately uniform in the direction of the positive  $\bar{x}$ -axis. This may be true, for example, if the body is almost plane and the oscillations are small. Then  $\Phi$  may be broken up into several parts, as

$$\Phi = U_\infty \bar{x} + \phi + \varphi \quad (2)$$

where the first term gives a uniform flow, the second term gives the correction for a steady flow about the body, the third term gives the correction to this for the oscillating body, and  $\phi$  and  $\varphi$  are small.

To the first order,  $\phi$  and  $\varphi$  are different solutions of the same differential equation

$$(1 - M^2) \varphi_{\bar{x}\bar{x}} + \varphi_{\bar{y}\bar{y}} + \varphi_{\bar{z}\bar{z}} - \frac{2 U_\infty}{a^2} \varphi_{\bar{x}t} - \frac{1}{a^2} \varphi_{tt} = 0 \quad (3)$$

where  $M, a$  are the Mach number and speed of sound at infinity. (See Reference 5, p. 198.)  $\varphi$  is a periodic function of  $t$ . Since the differential equation is linear, we may put  $\varphi = \bar{\varphi}(x, y, z) e^{i\omega t}$ , where  $\omega$  is the angular frequency of oscillation. In terms of the nondimensional quantities,

$$\begin{aligned}
x &= \bar{x}/b \\
y &= \bar{y}/b \\
z &= \bar{z}/b \\
k &= \omega b/U_\infty
\end{aligned}$$

(b is a characteristic length of the body); Equation (3) becomes

$$(1-M^2)\bar{\varphi}_{xx} + \bar{\varphi}_{yy} + \bar{\varphi}_{zz} - 2iM^2k\bar{\varphi}_x + M^2k^2\bar{\varphi} = 0 \quad (4)$$

For  $M = 1$ , this reduces to

$$\bar{\varphi}_{yy} + \bar{\varphi}_{zz} - 2ik\bar{\varphi}_x + k^2\bar{\varphi} = 0 \quad (5)$$

the linearized transonic equation (see Reference 6, p. 7). It has been suggested by Landahl (Reference 6) that the proper equation to use instead of (4) is

$$\bar{\varphi}_{yy} + \bar{\varphi}_{zz} - 2iM^2k\bar{\varphi}_x + M^2k^2\bar{\varphi} = 0$$

if  $k \gg |M-1|$ . Comparison of this equation with (5) leads to a similarity rule for flows in the transonic range (see Reference 6, p. 18).

The range of validity of this equation is discussed by Landahl (Reference 6, Chapter 1). First, there is the requirement for linearization in any speed range, that the perturbation potential  $\phi + \varphi$  be small. This is not satisfied at the leading edge of a wing for any realistic cross-sectional shape; however, it may be satisfied over the rest of the wing, if the wing has small thickness, and the results on parts of the wing away from the leading edge are not much affected by the error there.

Another restriction peculiar to transonic speeds is associated with the absence of the term in  $\bar{\varphi}_{xx}$ . The actual flow has some variation in local Mach number which may influence the nature of the flow considerably if  $M$  is near 1. The presence of the term in  $\bar{\varphi}_x$  tends to reduce this influence, but for  $k$  small or zero, the equation is valid only for a highly swept wing with a pointed nose.

The difference of the local Mach number from the value 1 assumed in Equation (5) may come from two sources: (1) wing thickness, and (2) a change in the free stream Mach number. Thus, for any value of  $k$ , there are limits on the thickness ratio and the Mach number range, which increase with  $k$ . Estimates of these limits are not possible, because of the small amount of experimental data available.

## 2. BOUNDARY CONDITIONS

The solution of Equation (1) must give a velocity field which is such that a particle at the body surface moves along the moving surface. If the equation of the surface is

$$S(\bar{x}, \bar{y}, \bar{z}, t) = 0$$

this equation must be satisfied when  $(\bar{x}, \bar{y}, \bar{z})$  moves with the velocity  $\nabla \Phi$ . Differentiating with respect to  $t$  gives the condition

$$\nabla \Phi \cdot \nabla S + \frac{\partial S}{\partial t} = 0 \quad (6)$$

This determines the normal velocity at the surface.

Now suppose that the body (to be referred to henceforth as a wing) is almost planar, lying almost in the  $xy$ -plane (see Figure 1). For vertical

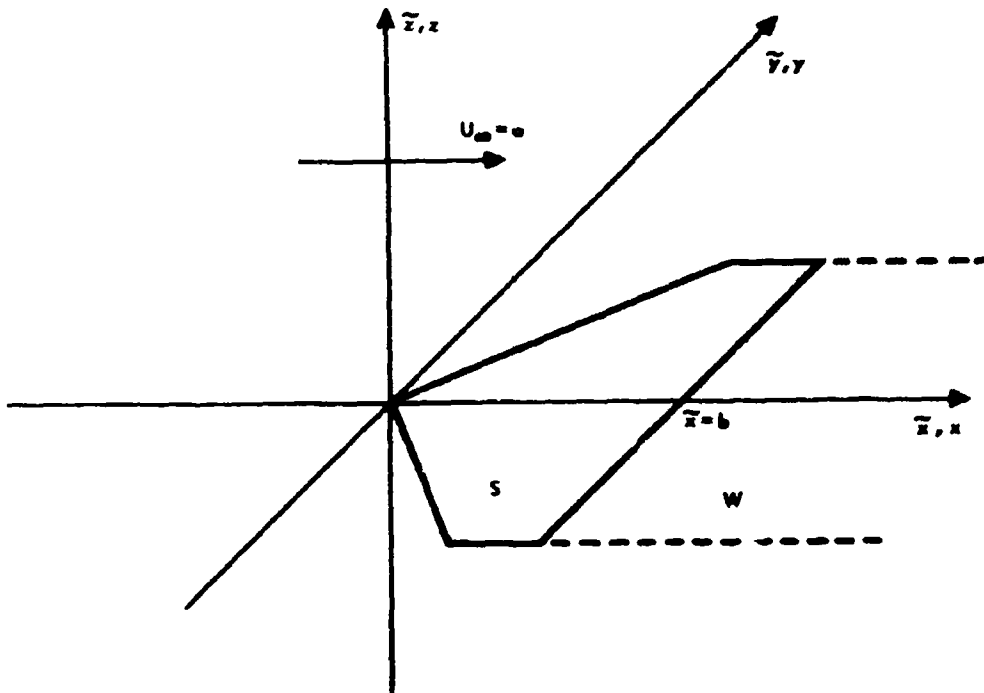


Figure 1. Coordinate Systems

oscillations of the body, the upper and lower surfaces may be represented by the equations

$$\bar{z} = \bar{g}_u(\bar{x}, \bar{y}) + e^{i\omega t} \bar{f}(\bar{x}, \bar{y})$$

$$\bar{z} = \bar{g}_u(\bar{x}, \bar{y}) + e^{i\omega t} \bar{f}(\bar{x}, \bar{y})$$

where the functions  $\bar{g}_u, \bar{g}_l$  are associated with the deviation of the shape of the body from planar, and  $\bar{f}$  depends on the mode of oscillation. Then on the two surfaces, we may take

$$S = \bar{z} - \bar{g}_u - e^{i\omega t} \bar{f}$$

$$S = \bar{z} - \bar{g}_l - e^{i\omega t} \bar{f}$$

Use these expressions for  $S$  and Equation (2) in Equation (6). Neglecting terms that involve products of  $\psi$  or  $\phi$  with  $\bar{g}_u, \bar{g}_l$ , or  $\bar{f}$ , the resulting equation may be broken up into a steady part, which gives the boundary condition for  $\phi$ , and an unsteady part, which gives the boundary condition for  $\bar{\varphi}$ . The unsteady part is

$$\frac{\partial \bar{\varphi}}{\partial z} = \frac{\partial \bar{f}}{\partial x} + ik\bar{f} \quad (7)$$

where  $f = \bar{f}/b$ . To the present degree of approximation, this condition should be applied at  $z = 0$ , over the region of the  $xy$ -plane on which the body projects.

### 3. THE BOUNDARY VALUE PROBLEM FOR $\bar{\varphi}$

In linearized theory, a disturbance of a flow at Mach 1 does not have any influence upstream. Consequently,

$$\bar{\varphi}(x, y, z) = 0, \quad x < 0 \quad (8)$$

if the body lies in the region  $x \geq 0$ . This is one of the conditions  $\bar{\varphi}$  must satisfy.

$\bar{\varphi}$  is a solution of Equation (5) in all space outside  $S$  and  $W$ , the regions in the  $xy$ -plane occupied by the wing and its wake (see Figure 1). In general,  $\bar{\varphi}$  is discontinuous in these regions. A boundary condition on  $W$  is obtained by equating the pressures above and below the surface of the wake. From the linearized form of the pressure coefficient,

$$\bar{C}_p = -2(\bar{\varphi}_x + ik\bar{\varphi})$$

(see Reference 6, p. 15) we get



$$\left[ \bar{\varphi}_x(x, y, z) + ik \bar{\varphi}(x, y, z) \right] \Big|_{z=0-}^{0+} = 0, \quad (x, y) \text{ in } W \quad (9)$$

This condition, plus Equation (7) applied on the two sides of  $S$ , plus Equation (8), determine  $\bar{\varphi}$  as a solution of Equation (5).

The conditions satisfied by  $\bar{\varphi}(x, y, z)$  are satisfied also by  $-\bar{\varphi}(x, y, -z)$ . Hence,  $\bar{\varphi}$  is an odd function of  $z$ . This implies that  $\bar{C}_p$  is zero in the wake. In the half space  $z > 0$ ,  $\bar{\varphi}$  is a solution of Equation (5), which satisfies Equation (8) and the boundary conditions

$$\bar{\varphi}_z(x, y, 0+) = \frac{\partial f}{\partial x} + ikf, \quad (x, y) \text{ in } S \quad (10)$$

$$\bar{\varphi}_x(x, y, 0+) + ik \bar{\varphi}(x, y, 0+) = 0, \quad (x, y) \text{ in } W \quad (11)$$

$$\bar{\varphi}(x, y, 0+) = 0, \quad (x, y) \text{ not in } S + W \quad (12)$$

Such a solution may be built up from a doublet distribution over  $S + W$  or a source distribution over the half plane  $z = 0$ ,  $x > 0$ .

#### 4. BASIC SOURCE AND DOUBLET SOLUTIONS OF THE DIFFERENTIAL EQUATION (See References 7, 8, and 9.)

The solution of Equation (5) which represents a point source at the origin is

$$\bar{\varphi}_0(x, y, z) = \begin{cases} 0, & x \leq 0 \\ -\frac{1}{2\pi} \frac{1}{x} e^{-\frac{1}{2} i k \left( x + \frac{y^2 + z^2}{x} \right)}, & x > 0 \end{cases} \quad (13)$$

(See Reference 9.) The potential of a point doublet oriented parallel to the  $z$ -axis is obtained by differentiation, as

$$\bar{\varphi}_1(x, y, z) = \frac{\partial \bar{\varphi}_0}{\partial z} = \begin{cases} 0, & x \leq 0 \\ \frac{i k}{2\pi} \frac{z}{x^2} e^{-\frac{1}{2} i k \left( x + \frac{y^2 + z^2}{x} \right)}, & x > 0 \end{cases} \quad (14)$$

It is easily verified that these functions satisfy Equation (5) for  $x \neq 0$ . They are poorly behaved at  $x = 0$  for real values of  $k$ .

To improve the behavior of  $\bar{\varphi}_0$  and  $\bar{\varphi}_1$  at  $x = 0$ , assume that  $k$  has a small negative imaginary part. This causes  $\bar{\varphi}_0$  and  $\bar{\varphi}_1$  to approach zero exponentially as  $x \rightarrow 0+$ , except at the origin. All partial derivatives of all orders have the same property. Thus,  $\bar{\varphi}_0$  and  $\bar{\varphi}_1$  are solutions of Equation (5) everywhere except at  $(0, 0, 0)$ . In the final formulas to be obtained, the imaginary part of  $k$  can be put equal to zero.

Solutions of Equation (5) for  $z > 0$  which satisfy Equation (8) are given for a distribution of sources as

$$\bar{\varphi}_s(x, y, z) = \iint_{\xi > 0} \rho(\xi, \eta) \bar{\varphi}_0(x-\xi, y-\eta, z) d\xi d\eta \quad (15)$$

and for a distribution of doublets as

$$\bar{\varphi}_d(x, y, z) = \iint_{\xi > 0} \rho(\xi, \eta) \bar{\varphi}_1(x-\xi, y-\eta, z) d\xi d\eta \quad (16)$$

where, to be completely general,  $\rho(\xi, \eta)$  may be any function such that the integrals exist. From the form of  $\bar{\varphi}_0$  and  $\bar{\varphi}_1$ , the region of integration may be restricted to the plane strip  $0 < \xi < x$ . It is shown in Appendix I that these functions satisfy the following boundary conditions for  $z = 0, x > 0$ :

$$\bar{\varphi}_{sz}(x, y, 0+) = \rho(x, y) \quad (17)$$

$$\bar{\varphi}_d(x, y, 0+) = \rho(x, y) \quad (18)$$

(in fact, if the same function  $\rho$  is used in both integrals,  $\bar{\varphi}_d = \partial \bar{\varphi}_s / \partial z$ ).

## 5. THE DETERMINATION OF $\bar{\varphi}$ BY A SOURCE DISTRIBUTION

One method of attack on the problem of finding  $\bar{\varphi}$  is to set  $\bar{\varphi} = \bar{\varphi}_s$ . Then, in terms of the upwash

$$w(x, y) = \bar{\varphi}_z(x, y, 0+) \quad (19)$$

we have from Equations (17) and (15)

$$\bar{\varphi}(x, y, z) = \iint_{\xi > 0} w(\xi, \eta) \bar{\varphi}_0(x-\xi, y-\eta, z) d\xi d\eta \quad (20)$$

for  $z \geq 0$ .

The values of  $w$  on  $S$  are known by Equation (10). Elsewhere,  $w$  is unknown, and it must be chosen so that the boundary conditions (11) and (12) are satisfied. We may take the limit as  $z \rightarrow 0+$  in Equation (20) by taking the limit under the integral sign:

$$\bar{\varphi}(x, y, 0+) = \iint_{\xi > 0} w(\xi, \eta) \bar{\varphi}_0(x-\xi, y-\eta, 0) d\xi d\eta \quad (21)$$

From Equations (11) and (12) are obtained the system of integral equations

$$\iint_{\xi > 0} w(\xi, \eta) \bar{\varphi}_0(x-\xi, y-\eta, 0) d\xi d\eta = 0, \quad (x, y) \text{ not in } S + W \quad (22)$$

$$\left( \frac{\partial}{\partial x} + ik \right) \iint_{\xi > 0} w(\xi, \eta) \bar{\varphi}_0(x-\xi, y-\eta, 0) d\xi d\eta = 0, \quad (x, y) \text{ in } W \quad (23)$$

Solution of Equations (22) and (23), followed by evaluation of  $\bar{\varphi}$  according to Equation (21), would yield the values of  $\bar{\varphi}$  on  $S$ , from which pressures and forces can be computed.

A box method based on a source distribution, described briefly in Reference 9, has been used by Weatherill at the Boeing Company. Some of his preliminary results are given in Reference 9.

## 6. THE DETERMINATION OF $\bar{\varphi}$ BY A DOUBLET DISTRIBUTION

If we set  $\bar{\varphi} = \bar{\varphi}_d$ , then by Equations (18), (16), and (12)

$$\bar{\varphi}(x, y, z) = \iint_{S+W} \bar{\varphi}(\xi, \eta, 0+) \bar{\varphi}_1(x-\xi, y-\eta, z) d\xi d\eta \quad (24)$$

In terms of

$$\psi(x, y) = \lim_{z \rightarrow 0} \frac{1}{z} \bar{\varphi}_1(x, y, z) = \begin{cases} 0, & x \leq 0 \\ \frac{ik}{2\pi} \frac{1}{x} e^{-\frac{1}{2}ik\left(x + \frac{y^2}{x}\right)}, & x > 0 \end{cases} \quad (25)$$

the normal derivative of  $\bar{\varphi}$  at  $z = 0$  is given by a singular integral:

$$w(x, y) = \iint_{S+W} \bar{\varphi}(\xi, \eta, 0+) \psi(x-\xi, y-\eta) d\xi d\eta \quad (26)$$

The values of  $\bar{\varphi}(\xi, \eta, 0+)$  must be determined then from

$$\iint_{S+W} \bar{\varphi}(\xi, \eta, 0+) \psi(x-\xi, y-\eta) d\xi d\eta = w(x, y), \quad (x, y) \text{ in } S \quad (27)$$

$$\left( \frac{\partial}{\partial x} + ik \right) \bar{\varphi}(x, y, 0+) = 0, \quad (x, y) \text{ in } W \quad (28)$$

## 7. A COMPARISON OF THE METHODS

Except for the singularity of the integral in Equation (27), all points of difference are in favor of solving the problem by doublets. There are these points:

- a. The region of integration in the source method extends theoretically to  $\pm\infty$  in  $\eta$ ; even practically, the region must be extended an extreme distance. In the doublet method, the region is restricted to  $S + W$ . This distinction is not so great for supersonic flows. There, the region of influence of the wing is swept back along Mach lines, and the set of points in this region that influences the wing is bounded (see Reference 10).
- b. After the unknown function under the integral sign is known, the source method requires an extra step – the evaluation of  $\bar{\varphi}$  on the wing from Equation (21).
- c. If values in the wake must be considered, the condition in the wake for the source method, Equation (23), is more complicated than the corresponding condition, Equation (28), for the doublet method.

The doublet method was used because of point a.

## 8. THE ADVANTAGE OF A STRAIGHT TRAILING EDGE

Suppose the wing has a straight trailing edge perpendicular to the direction of flow ( $x = \text{constant}$  along the edge); then the wing is not influenced by the wake. This is reflected in the equations by the fact that the integrands are zero when  $\xi > x$ . Hence, in either method, for the determination of  $\bar{\varphi}$  on the wing, the condition in  $W$  need not be used.

## 9. THE DOUBLET BOX METHOD

Consider a flow at Mach 1 past an oscillating wing with its nose at the origin, lying approximately in the  $xy$ -plane, with  $x = l$  along the trailing edge. The value of the unsteady potential  $\bar{\varphi}$  on the wing may be found by solution of Equation (27), which may be written as

$$\iint_S \bar{\varphi}(\xi, \eta, 0+) \psi(x-\xi, y-\eta) d\xi d\eta = w(x, y), \quad (x, y) \text{ in } S \quad (29)$$

To get an approximate solution of this equation, let the  $xy$ -plane be covered with a grid of square boxes with sides of length  $d$ , so that box edges lie along the coordinate axes (see Figure 2). Let the region  $B$  be composed of all boxes whose centers lie in  $S$ ;  $B$  is an approximation to  $S$  by boxes. Let  $i, j$  be box indexes in the  $x$ - and  $y$ -directions. Approximate  $\bar{\varphi}$  by a constant value  $\bar{\varphi}_{ij}$  in the  $(i, j)$ -th box  $B_{ij}$ . Impose the condition of Equation (7) at the center  $(x_i, y_j)$  of each box  $B_{ij}$  in  $B$ , with the region of integration replaced by  $B$ . Then Equation (29) gives a system of linear algebraic equations for the  $\bar{\varphi}_{ij}$ 's:

$$\sum_{i', j'} \bar{\varphi}_{i'j'} \iint_{B_{i'j'}} \psi(x_i - \xi, y_j - \eta) d\xi d\eta = w(x_i, y_j) \quad (30)$$

Examination of the integral in Equation (30) shows that it depends on  $i, j, i', j'$  only via  $i-i', |j-j'|$ . The notation

$$A(i-i', |j-j'|) = \iint_{B_{i'j'}} \psi(x_i - \xi, y_j - \eta) d\xi d\eta \quad (31)$$

is introduced. Formulas for the evaluation of this quantity are given in Appendix II.

Segregating the terms with  $i'=i$  on the left, Equation (30) becomes

$$\sum_{j'} A(0, |j-j'|) \bar{\varphi}_{ij'} = w(x_i, y_j) - \sum_{i' < i} \sum_{j'} A(i-i', |j-j'|) \bar{\varphi}_{i'j'} \quad (32)$$

For fixed  $i$  and varying  $j$ , this is a smaller system of equations that may be solved for each consecutive value of  $i$ .

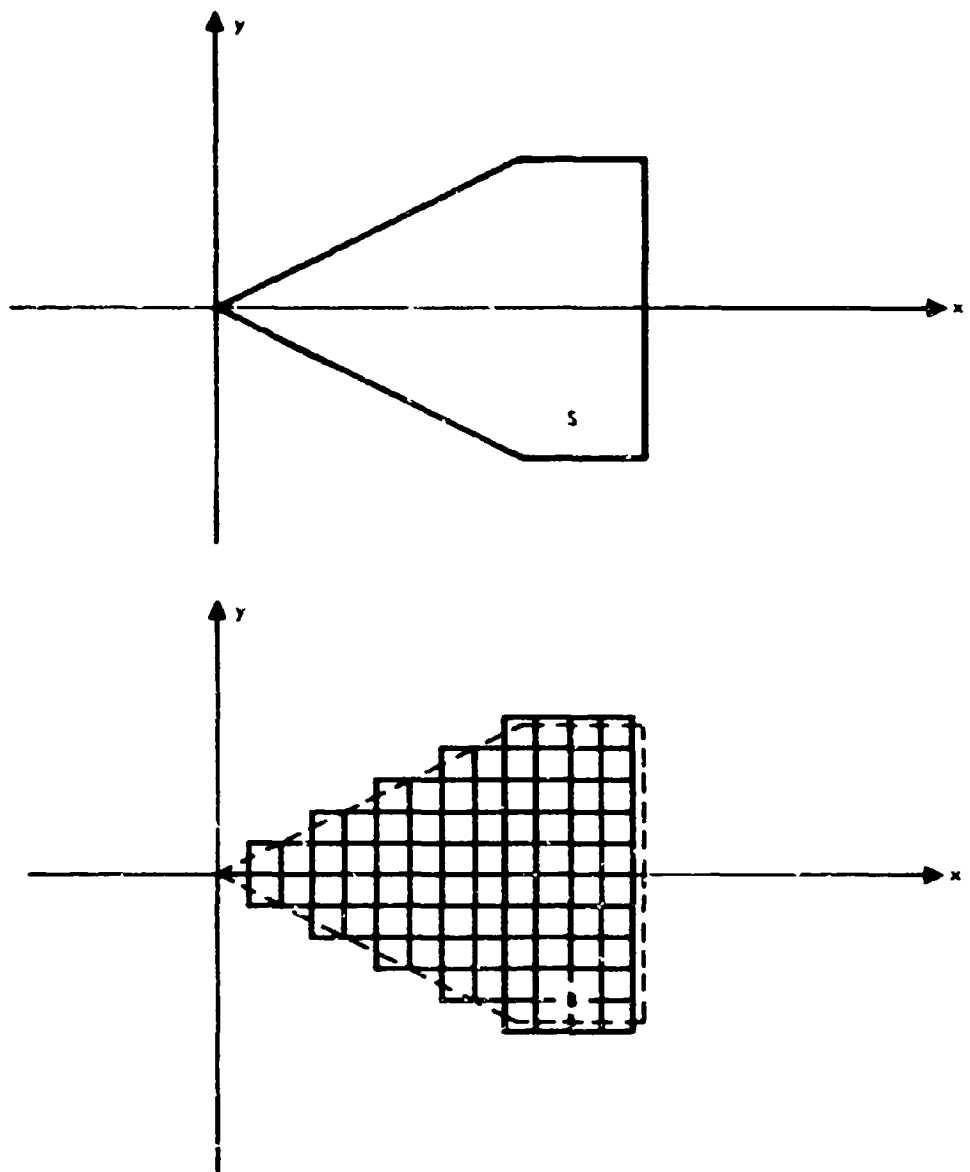


Figure 2. Approximation of the Wing by Region B

Now suppose the wing is symmetric about the x-axis; then only modes of oscillation that are symmetric or antisymmetric in y need be treated. Consider a symmetric mode.  $\bar{\varphi}_{ij}$  will have the same value at corresponding boxes across the x-axis. This may be used to reduce the range of the sum in Equation (32) and the range of j. Let j = 1 in the row of boxes in which  $0 < y < d$ . Then, combining terms for symmetrically placed boxes,

$$\sum_{j' \geq 1} [A(0, |j-j'|) + A(0, j+j'-1)] \bar{\varphi}_{ij'} \quad (33)$$

$$= w(x_i, y_j) - \sum_{i' < i} \sum_{j' \geq 1} [A(i-i', |j-j'|) + A(i-i', j+j'-1)] \bar{\varphi}_{i'j'}, \quad j \geq 1$$

The equations for  $j \leq 0$  are implied by those with  $j \geq 1$ . Thus, the size of the system has been reduced by a factor of 2.

For antisymmetric modes, Equation (33) applies, with the sums of values of A replaced by differences.

## 10. EXTENSIONS OF THE METHOD

The computer program discussed in Section 3 has some restrictions that are not inherent in the box method, such as the requirement of a straight trailing edge. Some possible modifications that extend the applicability of the program will now be described.

To modify the program for modes antisymmetric in y, it is only necessary to change some of the signs in Equation (33), as indicated in the discussion above, and replace even powers of y by odd powers in the formulas used for deflection and potential.

To deal with a more general trailing edge, it is necessary to use the values of  $\bar{\psi}$  in the wake. For fixed y, if  $x = x_T$  at the trailing edge, Equation (28) may be integrated to give

$$\bar{\psi}(x, y, 0+) = e^{-ik(x - x_T)} (x_T, y, 0+)$$

in W. In addition to the set of boxes B on the wing, a corresponding set of boxes  $B_W$  on W must be considered. After finding a value  $\bar{\psi}_{ij}$  in a box of B along the trailing edge, the formula above may be used to find values in the boxes directly downstream. If the ith row of boxes includes boxes of  $B_W$ , to the right side of Equation (33) must be added the contribution of all boxes  $B_{i'j'}$  in  $B_W$  with  $i' \leq i$ . The computer program must also be modified in several other respects, to take into account the more general wing shape.

A wing that consists of several almost planar sections in different planes, such as a wing with folded tips, may also be handled by the doublet box method. Equation (33) applies, if  $\bar{\psi}_{ij}$  is interpreted as one-half of the discontinuity in  $\bar{\psi}$  between the upper and lower surfaces. The influence coefficients involved are given by a more general formula (not given in this report), allowing for out-of-plane influence of the doublets. Formulas analogous to those of Appendix II may be developed, which are not much more complicated. The main effect of this extension on the computer program would be a greater number of distinct values of the influence coefficients, so that it would not be possible to store them all in an array in core unless the limit on the number of boxes in each direction were considerably reduced.

Rectangular boxes, not necessarily square, may be used. Let the boxes have sides of length  $d_1$  chordwise and  $d_2$  spanwise.

If

$$l_1 = kd_1, l_2 = kd_2^2/d_1$$

the formula for the influence coefficients, Equation (39) in Appendix II, must be replaced by

$$A(n, m) = \frac{ik}{2\pi} \iint_{\substack{|v-m| < 1/2 \\ |u-n| < 1/2 \\ u > 0}} \frac{dudv}{u^2} e^{-1/2 i (l_1 u + l_2 v^2/u)}$$

This may be evaluated by the methods of Appendix II. Except for this difference, the method is essentially the same. The best choice of box shape probably depends on the aspect ratio of the wing.



### 3. DESCRIPTION OF THE COMPUTER PROGRAM

#### 1. COORDINATE SYSTEMS

An initial coordinate system  $(\bar{x}, \bar{y}, \bar{z})$  is assumed, with the  $\bar{x}$ -axis in the direction of the flow. The undisturbed position of the wing is in a region  $S$  in the  $\bar{x}\bar{y}$ -plane, with the  $\bar{x}$ -axis along the center line and the origin at the nose (see Figure 1). This coordinate system is used in the data.

In the program, a dimensionless coordinate system  $(x, y, z)$  is used, based on the root chord length  $b$ :

$$x = \bar{x}/b$$

$$y = \bar{y}/b$$

$$z = \bar{z}/b$$

#### 2. WING GEOMETRY

The wing is symmetric, with trailing edge  $\bar{x} = b$ . To complete its description, the portion of the leading edge on which  $\bar{y} > 0$  must be specified. This is done by giving the coordinates of the end points of  $NS$  line segments along the edge ( $1 \leq NS \leq 3$ ), beginning at a point at which  $\bar{y} = 0$ :  $(0, \bar{y}_0)$ ,  $(\bar{x}_1, \bar{y}_1)$ , . . . ,  $(\bar{x}_{NS}, \bar{y}_{NS})$ . The edge of  $S$  includes the polygonal line through these points. If  $\bar{y}_0 > 0$ , it also includes the line from the origin to  $(0, \bar{y}_0)$ . If  $\bar{x}_{NS} < b$ , it includes the line from  $(\bar{x}_{NS}, \bar{y}_{NS})$  to  $(b, \bar{y}_{NS})$ . (See Figure 3.)

Leading edges of fairly general shape may be approximated by such polygonal lines.

#### 3. THE DEFLECTION DATA

A mode is specified by the vertical deflection function  $f(\bar{x}, \bar{y})$  in terms of which the equation for the instantaneous position of the planform is

$$\bar{z} = \text{Re} [\delta \cdot e^{i\omega t} f(\bar{x}, \bar{y})]$$

where  $\delta$  is a constant.

In the program,  $f$  is assumed to be a polynomial in  $\bar{x}$  and  $\bar{y}^2$ . The data may give either the coefficients of this polynomial, or values of  $f$  at a

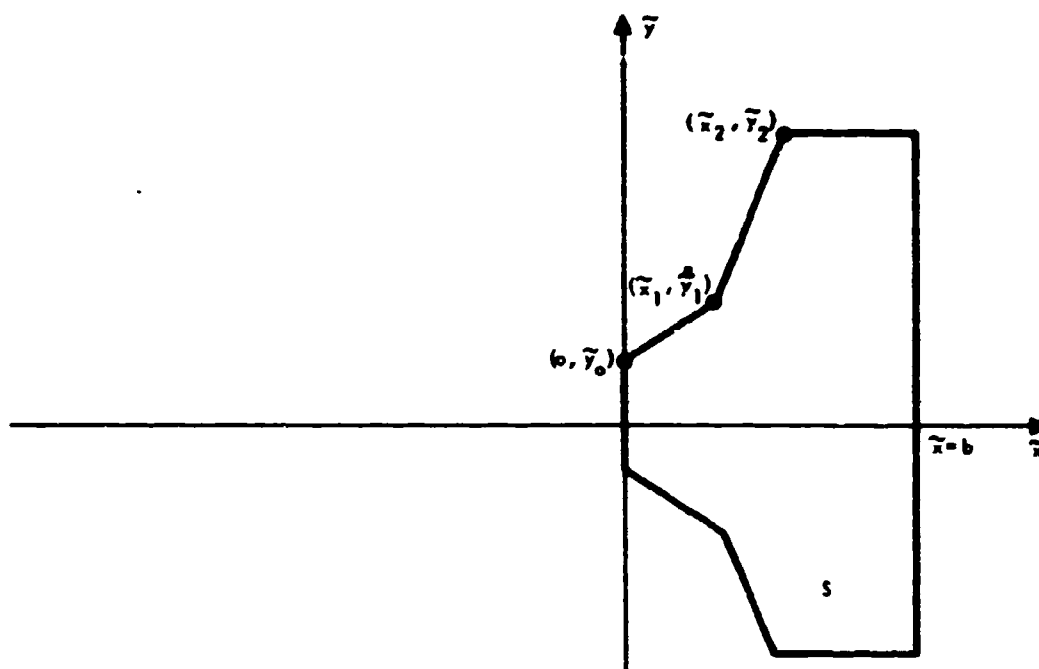


Figure 3. Wing Geometry (NS = 2)

set of points on the wing. In the latter case, a polynomial is fitted to the given values by a least square error technique.

#### 4. LEAST SQUARE SURFACE FITS

The problem involved here is the approximation of a function of  $x$  and  $y$  by an expression of the form

$$\bar{\varphi}(x, y) = \sum_{n, m} a_{nm} x^n y^{2m} F(x, y)$$

when a set of values of the function is known. This arises in the program in two places. The subroutine DRED fits a representation of the deflection of this type with  $F = 1$  to the given deflection values. In the subroutine B0XP, such a fit is made for the potential, with

$$F(x, y) = \left\{ \begin{array}{l} \sqrt{x^2 - x_0(y)^2} \\ \text{or} \\ \sqrt{x - x_0(y)} \end{array} \right\} \cdot \left\{ \begin{array}{l} \sqrt{1 - y^2/y_{\max}^2} \\ \text{or} \\ 1 \end{array} \right\}$$

( $x = x_0(y)$  is the equation of the leading edge) depending on the wing shape. This factor approximates the proper behavior of  $\bar{\varphi}$  at the edges. The factor  $\sqrt{x^2 - x_0(y)^2}$  is used for a pointed nose ( $\tilde{y}_0 = 0$ ), and  $\sqrt{x - x_0(y)}$  for an unswept nose ( $\tilde{y}_0 > 0$ ). The factor  $\sqrt{1 - y^2/y_{\max}^2}$  is included if the planform has a side edge along which  $y = y_{\max}$ .

The factor  $F(x, y)$  is real, so the values of  $\bar{\varphi}$  have real and imaginary parts that involve only the corresponding parts of the  $a_{nm}$ 's. Hence, these real and imaginary parts may be handled separately, reducing the problem from one in complex numbers to one in real numbers.

Let  $\alpha_{nm} = \text{Re}[a_{nm}]$ , and let the real parts of given values of the function at data points be  $\bar{\varphi}'_j$  at  $(x_j, y_j)$ ,  $j = 1, \dots, NP$ . Then for the real parts we wish to have

$$\sum_{n, m} \alpha_{nm} x_j^n y_j^{2m} F(x_j, y_j) \cong \bar{\varphi}'_j, \quad j = 1, \dots, NP$$

The least squares method minimizes

$$Q = \sum_j \left[ \sum_{n, m} \alpha_{nm} x_j^n y_j^{2m} F(x_j, y_j) - \bar{\varphi}'_j \right]^2$$

(See Reference 11, Chapter 16.)

For condensed notation, let  $r$  be a single index over the pairs  $(n, m)$ , let  $\alpha_{nm} = \alpha_r$ , and  $x_j^n y_j^{2m} F(x_j, y_j) = A_{jr}$ . Then

$$Q = \sum_j \left[ \sum_r \alpha_r A_{jr} - \bar{\varphi}'_j \right]^2$$

Let the range of  $r$  be from 1 to  $NC \leq NP$ .

To minimize  $Q$ , we set

$$\frac{\partial Q}{\partial \alpha_r} = 0, \quad r = 1, \dots, NC$$

This leads to the system of equations

$$\sum_{r'} \left( \sum_j A_{jr} A_{jr'} \right) \alpha_{r'} = \sum_j A_{jr} \bar{\varphi}_j', \quad r = 1, \dots, NC \quad (34)$$

Put

$$\left. \begin{aligned} \sum_j A_{jr} A_{jr'} &= B_{rr'} \\ \sum_j A_{jr} \bar{\varphi}_j' &= C_r' \end{aligned} \right\} \quad (35)$$

Then Equation (34) reduces to

$$\sum_{r'} B_{rr'} \alpha_{r'} = C_r', \quad r = 1, \dots, NC \quad (36)$$

The matrices  $(B_{rr'})$  and  $(C_r')$  must be set up to solve Equation (36). It is not necessary, however, to set up the matrix  $(A_{jr})$ . Only one row of  $(A_{jr})$  is needed at a time. This is fortunate, because the program allows  $(A_{jr})$  to become as large as  $2500 \times 20$ . For each value of  $j$ , the  $j$ th row of  $(A_{jr})$  is computed, and from this the  $j$ th terms in the sums in Equation (35) are formed and added in.

In the complex case, there is a corresponding system of equations for the imaginary parts:

$$\sum_{r'} B_{rr'} \beta_{r'} = C_r''$$

The two systems of equations are solved together by the subroutine XSIMEQ, which allows for more than one set of values on the right.

## 5. GENERALIZED FORCES

The generalized force coefficient  $L_{ij}$  is defined (Reference 6) by

$$L_{ij} = \frac{1}{1/2 \rho U_\infty^2 S} \iint_S \Delta p_i(x, y) f_j(x, y) dx dy$$

where  $\Delta p_i$  is the lifting pressure difference in the  $i$ th mode, and  $f_j$  is the deflection function in the  $j$ th mode. In terms of the potential  $\bar{\varphi}(x, y)$  on the upper surface,

$$\Delta p_i = 2\rho U_\infty^2 (\bar{\varphi}_x + i k \bar{\varphi})$$

$$L_{ij} = \frac{4}{S} \iint_S (\bar{\varphi}_x + i k \bar{\varphi}) f_j dx dy$$

After integration by parts,

$$L_{ij} = \frac{4}{S} \left\{ \int_{x=1} \bar{\varphi} f_j dy + \iint_S \bar{\varphi} \left( i k f_j - \frac{\partial f_j}{\partial x} \right) dx dy \right\} \quad (37)$$

In Equation (37) insert the series

$$\bar{\varphi} = \sum_{n, m} a_{nm} x^n y^{2m} F(x, y)$$

$$f_j = \sum_{n', m'} d_{n'm'} x^{n'} y^{2m'}$$

The result is

$$\begin{aligned} L_{ij} = \frac{8}{S} \sum_{n', m'} d_{n'm'} \sum_{n, m} a_{nm} & \left[ \frac{1}{2} \int_{x=1} y^{2m+2m'} F(1, y) dy \right. \\ & + i k \cdot \frac{1}{2} \iint_S x^{n+n'} y^{2m+2m'} F(x, y) dx dy \\ & \left. - n' \cdot \frac{1}{2} \iint_S x^{n+n'-1} y^{2m+2m'} F(x, y) dx dy \right] \end{aligned}$$

The integrals in this expression depend only on the wing shape. They are computed by the subroutine FØRCI before the work on the individual modes begins. During the work on the  $i$ th mode, the sum over  $n$  and  $m$  is

performed in the last part of the subroutine B0XP, for each set of values of  $n'$  and  $m'$ . The sum over  $n'$  and  $m'$  and multiplication by  $8/S$  is performed in the last part of the main program.

## 6. THE USE OF GAUSSIAN QUADRATURE IN THE EVALUATION OF GENERALIZED FORCES

Gaussian quadrature is an approximation of the form

$$\int_a^b f(u) du \cong \sum_{j=1}^N h_j f(u_j)$$

exact for polynomials of degree  $\leq 2N - 1$ . (See Reference 11, Chapter 7.) This formula is used with  $(a, b) = (0, 1)$ ,  $N = 6$ . The values of the  $h_j$ 's and  $u_j$ 's for this case were obtained from values listed in Reference 11. They are given as  $H(1), \dots, H(6)$ ,  $U(1), \dots, U(6)$  in the subroutine SECT.

Subroutine F0RCI finds the values of

$$AXY(I, J) = \frac{1}{2} \iint_S x^{I-1} y^{2J-2} F(x, y) dx dy$$

and

$$AY(J) = \frac{1}{2} \int_{x=1} y^{2J-2} F(1, y) dy$$

for  $I, J = 1, \dots, 9$ . To do this, the contributions to the integrals from each section of wing behind a straight piece of leading edge are calculated separately in SECT.

The form of  $F(x, y)$  is

$$F(x, y) = \left\{ \begin{array}{l} \sqrt{x - x_0(y)} \\ \text{or} \\ \sqrt{x^2 - x_0(y)^2} \end{array} \right\} \cdot \left\{ \begin{array}{l} \sqrt{1 - y^2/y_{\max}^2} \\ \text{or} \\ 1 \end{array} \right\}$$

depending on the wing shape. We have integrals that behave like square roots at the leading edge. The integrals over one wing section are of the form

$$BXY = \int_{y_-}^{y_+} dy \int_{x_0(y)}^1 dx x^{I-1} y^{2J-2} F(x, y)$$

and

$$BY = \int_{y_-}^{y_+} dy y^{2J-2} F(1, y)$$

In BXY, the chordwise integral is evaluated first at each value of  $y$  at which it will be needed. The new variable

$$u = \sqrt{x - x_0(y)} / \sqrt{1 - x_0(y)} \quad (38)$$

is introduced. Then

$$\int_{x_0(y)}^1 dx x^{I-1} y^{2J-2} F(x, y) = \int_0^1 du \cdot 2 \left[ 1 - x_0(y) \right] x^{I-1} y^{2J-2} F(x, y)$$

The integrand, as a function of  $u$ , is well-behaved at the leading edge. It is approximated by

$$\sigma(y) = \sum_{i=1}^6 h_i \cdot 2 \left[ 1 - x_0(y) \right] x_i^{I-1} y^{2J} F(x_i, y)$$

where  $x_i$  is computed from the value of  $u_i$  according to Equation (38).

In the  $y$ -integration in BXY and BY, the integrand approaches zero as  $y \rightarrow y_{\max}$  like  $\sqrt{1 - y/y_{\max}}$  or  $(1 - y/y_{\max})^{3/2}$ . Accordingly, the change of variable

$$y = \begin{cases} y_+ - (y_+ - y_-)v, & y_+ < y_{\max} \\ y_+ - (y_+ - y_-)v^2, & y_+ = y_{\max} \end{cases}$$

is used, which makes the interval of integration  $0 < v < 1$  and removes the square root behavior in the last section of the wing. This leads to the formulas

$$\begin{aligned}
 BXY &= (y_+ - y_-) \sum_{j=1}^6 h_j \sigma(y_j) \cdot \begin{cases} 1, y_+ < y_{\max} \\ 2u_j, y_+ = y_{\max} \end{cases} \\
 BY &= (y_+ - y_-) \sum_{j=1}^6 h_j y_j^{2J-2} F(1, y_j) \cdot \begin{cases} 1, y_+ < y_{\max} \\ 2u_j, y_+ = y_{\max} \end{cases}
 \end{aligned}$$

## 7. LEADING EDGE CORRECTION

The value of potential found for each box from Equation (33) is taken to be the value of  $\bar{\varphi}$  at the box center. Thus, the values obtained are in error only by virtue of the error introduced in the values of upwash when the actual distribution of potential in a box is replaced by this constant value. This error is especially important in the first row of boxes, for a wing with an unswept leading edge. The major effect is on the upwash values in that row.

To estimate this error, consider the two-dimensional case, in which  $\bar{\varphi}$  is independent of  $y$ . In Equation (26), the expression for upwash due to a doublet distribution, integrate by parts over  $\xi$ , then integrate over  $\eta$ . The result is

$$\begin{aligned}
 \bar{w}(x, y) &= \frac{ik}{2\pi} \iint_{0 < \xi < x} \frac{d\xi d\eta}{(x - \xi)^2} \bar{\varphi}(\xi, \eta, 0+) e^{-\frac{1}{2} ik \left( x - \xi + \frac{(y - \eta)^2}{x - \xi} \right)} \\
 &= \frac{1}{\pi} \iint_{0 < \xi < x} \frac{d\xi d\eta}{(y - \eta)^2} \left( \bar{\varphi}_\xi + \frac{1}{2} ik \bar{\varphi} \right) e^{-\frac{1}{2} ik \left( x - \xi + \frac{(y - \eta)^2}{x - \xi} \right)} \\
 &= -\sqrt{\frac{2ik}{\pi}} \int_0^x \frac{d\xi}{\sqrt{x - \xi}} e^{-\frac{1}{2} ik(x - \xi)} \left( \bar{\varphi}_\xi + \frac{1}{2} ik \bar{\varphi} \right)
 \end{aligned}$$

For  $x = \frac{1}{2} d$ , if  $kd$  is small,



$$\bar{w}\left(\frac{1}{2}d, y\right) \cong -\sqrt{\frac{2ik}{\pi}} \int_0^{\frac{1}{2}d} \frac{d\xi}{\sqrt{\frac{1}{2}d - \xi}} \bar{\varphi}_{\xi}(\xi, \eta, 0+)$$

The correct leading edge behavior is possessed by the expression  $\bar{\varphi} = C\sqrt{\xi}$ . We have

$$\bar{w}\left(\frac{1}{2}d, y\right) \Big|_{\bar{\varphi} = C\sqrt{\xi}} = -\sqrt{\frac{2ik}{\pi}} \frac{C}{2} \int_0^{\frac{1}{2}d} \frac{d\xi}{\sqrt{\xi\left(\frac{1}{2}d - \xi\right)}} = -\sqrt{\frac{2ik}{\pi}} C \frac{\pi}{2}.$$

If  $\bar{\varphi}$  is constant on  $0 < \xi < d$ , and has the value  $C\sqrt{1/2}d$ , then  $\bar{\varphi}_{\xi}$ , in the above integral, can be expressed in terms of a delta function:

$$\bar{\varphi}_{\xi} = C\sqrt{\frac{1}{2}d} \delta(\xi)$$

Accordingly,

$$\bar{w}\left(\frac{1}{2}d, y\right) \Big|_{\bar{\varphi} = C\sqrt{\frac{1}{2}d}} = -\sqrt{\frac{2ik}{\pi}} C.$$

Note that the latter value is smaller than the value of  $\bar{w}$  evaluated for  $\bar{\varphi} = C\sqrt{\xi}$  by the factor  $2/\pi$ . This implies that the values of potential found for the first row of boxes will be more accurate if the upwashes in that row are multiplied by  $2/\pi$ .

## 8. THE FORM OF OUTPUT

The viewpoint is taken that calculation of the generalized forces is the basic purpose of the program. They are always printed out. There are other outputs that will be printed if the appropriate data signal is given. Each of the following is printed if the data item specified in parentheses is non-zero:

- a. The coefficients of the deflection polynomial, if it has been computed as a fit to given values of deflection, DA(87)
- b. The upwash array, DA(88)
- c. The potential array, DA(89)
- d. The coefficients of the potential series, DA(90)

- e. Values of pressure and potential at the box centers, computed from the series, DA(91).

## 9. THE DATA SUBROUTINE DATRD

This subroutine reads all data items into the array DA. Punched cards used for data are considered to contain six fields of length 12 as indicated in the sample data sheets. The first field contains information for DATRD. Ending in column 12 is an integer giving the location in the data array for the entry in the second field. The following fields go into consecutive locations, if the data are numeric. Floating point numbers should be written with decimal points, and fixed point numbers adjusted to the right end of the field.

The word ALPHA in columns 2 through 6 indicates that the data on the card are alphanumeric. These are stored in DA in a different way, taking up ten locations per card. The data may be printed later, just as they appeared on the card.

On a numeric card, if a field is blank the corresponding location in DA is unchanged. This is not true for an ALPHA card.

A minus sign in column 1 indicates the last card to be read at the time. DATRD reads cards until this minus sign is encountered, then returns to the main program.

## 10. A NOTE ON THE USE OF TAPES

In writing of this program, the following tape numbers have been used: output tape, number 6; input tape, number 5; and tape simulated by an internal file, number 99.

The tape numbers 5 and 99 appear in the subroutine DATRD. Elsewhere only the output tape number is used. It occurs in the main program, and in the subroutines SHAPE, DRED, B0XP and B0XP0.

## 11. USE OF THE PROGRAM FOR FIXED WING AND MODES AT VARIOUS FREQUENCIES

If a non-zero quantity is entered in the appropriate location in the data array, (DA26), it indicates that a wing shape and set of modes to be used are the same already used for another frequency. Then quantities that depend only on wing shape and deflection data will not be computed, but will be taken from the permanent arrays in which they were stored in the previous case. The number of boxes along the root chord, DA(27), may not be changed when this is done.

When this option is exercised, all work for the present frequency will be carried out after reading one set of data, which need only include the frequency and the indicator, DA (26). Titles for the individual modes are not printed.

## 12. DESCRIPTION OF THE DATA ARRAY

All data are entered into the array DA, dimensioned for 700, as described in Paragraph 9. The layout of the array is as follows:

1 - 10	Title
13 - 22	Mode Title
23 :	Frequency (cycles per unit of time), $\nu$
24 :	Root chord length, $b$
25 :	Speed of sound, $a$
26 :	Indicator for new frequency (See Paragraph 11)
27 :	Number of boxes along root chord
28 :	Number of modes
29 :	Number of sections of leading edge to be given (See Paragraph 2)
30 - 36	Coordinates of points on the leading edge (See Paragraph 2)
39 :	Indicator to suppress calculation of potential for a mode
46 - 70	Coefficients of the deflection polynomial (See Paragraph 3)
87 - 91	Output indicators (See Paragraph 8)
98 :	Number of points at which deflections are given (See Paragraph 3)
99 :	Number of $\bar{x}$ values
100 :	Number of $\bar{y}$ values
101 - 700	Deflection data for a maximum of 150 points

Note: 23, 24, 25 and 30-36 must be entered in consistent units of length and time.

### 13. OUTLINE OF THE PROGRAM

For the purpose of description, the main program has been divided into 20 parts, as indicated in Figure 4, which shows the flow of the program and the subroutines called.

### 14. SIZE LIMITATIONS OF THE PROGRAM

The program's size limitations are as follows:

- a. Box size - the half wing must be enclosed in a rectangle that contains no more than 50 boxes in each direction. (The use of a large number of boxes is not recommended, because the time required is roughly proportional to the cube of the number of boxes along the root chord. The possibility of 50 boxes in each direction is intended to allow a large range in aspect ratio.)
- b. Number of modes - ten at most.
- c. Number of points at which the deflections are given for one mode - 150 at most.
- d. Terms in the deflection polynomial - this is  $\sum_{nm} d_{nm} x^n y^{2m}$ , where  $0 \leq n \leq 4$ ,  $0 \leq m \leq 4$ .

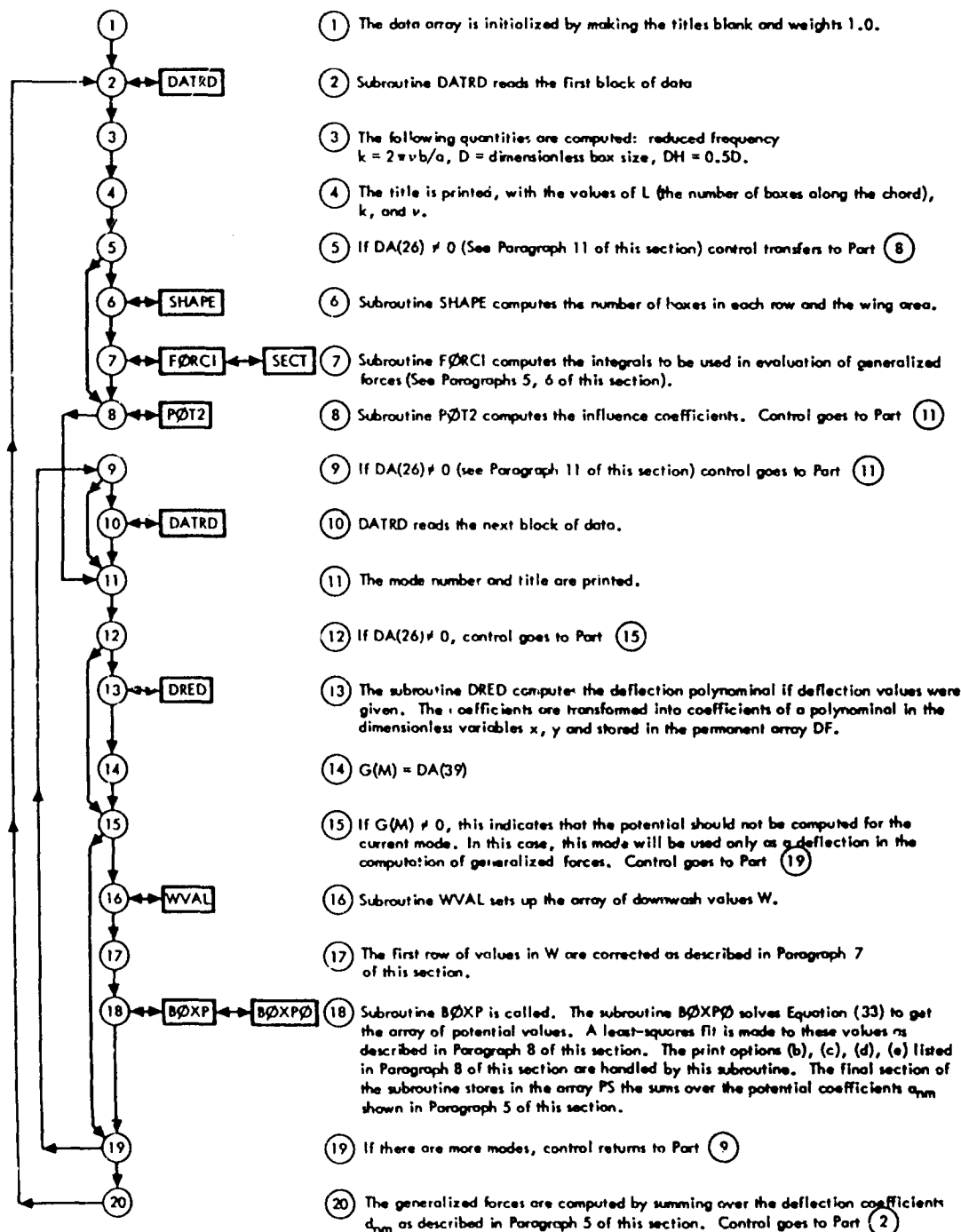


Figure 4. Program Flow Diagram

## 4. RESULTS

### 1. THE ASPECT RATIO 1.5 DELTA WING

The computer program was run for the plunging and pitching modes (pitch axis at  $x = 0$ ) at the reduced frequencies  $k = 0.2, 0.5, 0.8, 1.0$ . Forty boxes along the root chord were used, which leads to about 300 boxes on the half-wing.

Theoretical values for comparison were calculated from Davies' formulas (Reference 12). These are analytic expressions of the solution of Equation (5) for the potential and generalized forces for the delta wing in rigid modes of oscillation, expressed as series in  $k$ . Figures 5 through 7 show the values of generalized forces  $L_{11}$  (lift due to plunge),  $L_{21}$  (lift due to pitch), and  $L_{22}$  (moment due to pitch). Note that the vertical scales have been expanded in the portion of interest, especially for  $L_{11}$ . Most of the values agree to within 2 or 3 percent.

The differences indicate the errors introduced by the box method in the solution of Equation (5), as distinguished from the errors inherent in this equation.

Figure 8 gives the chordwise distribution of values of  $\bar{\phi}$  for the plunging mode at  $k = 0.5$ , for  $y = y_{\max}/3 = 0.125$ .

### 2. THE ASPECT RATIO 2.0 RECTANGULAR WING

The plunging and pitching modes were again used at  $k = 0.3, 0.6, 0.9$ . Twenty-five boxes were allowed along the chord, giving 625 boxes on the half-wing. The values of  $L_{21}$  and  $L_{22}$  are shown in Figures 9 and 10, with values from Landahl (Reference 6, page 84) for comparison. Landahl's values were obtained by a method of solution of Equation (5) which applies only to a rectangular wing in modes of oscillation with a deflection independent of  $y$ .

### 3. THE ASPECT RATIO 3.0 RECTANGULAR WING

Finally, for the aspect ratio 3.0 rectangular wing, a comparison is made with experimental pressure values. These values were given in Reference 13 for a 5-percent thickness wing oscillating in an elastic bending mode. At Mach 1, the reduced frequency was 0.24. The chordwise pressure distribution at  $y = y_{\max}/2$  is shown in Figure 11.

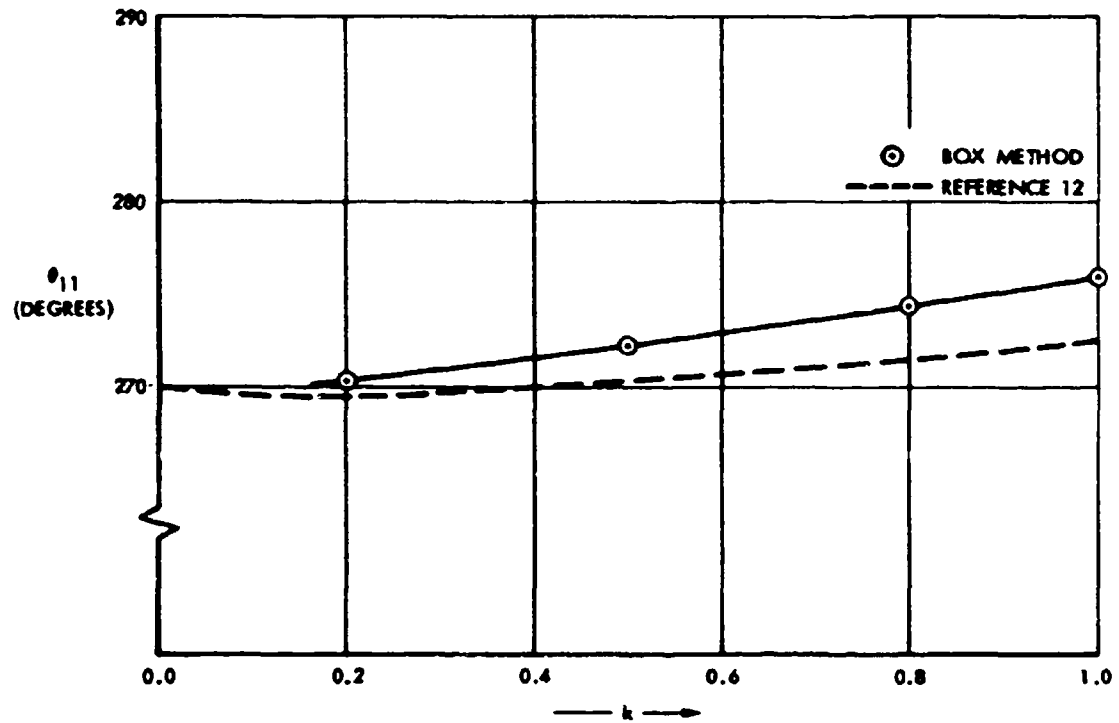
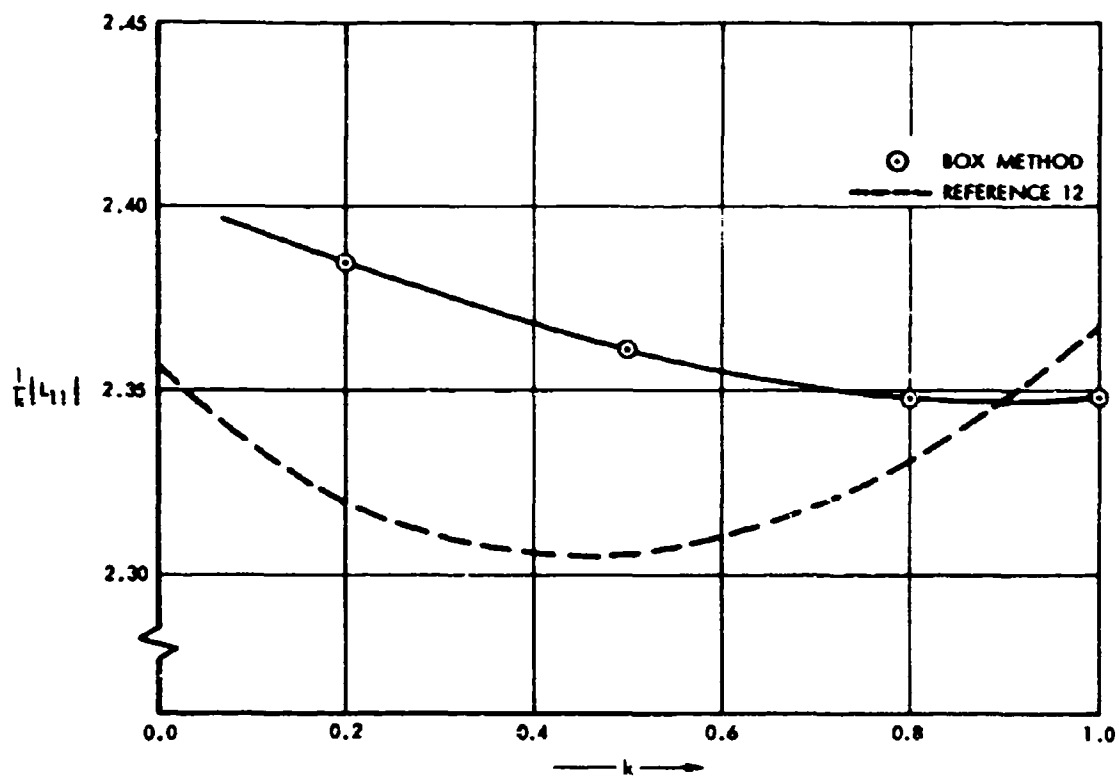


Figure 5. Lift Due to Translation for an Aspect Ratio 1.5 Delta Win  
(Compared With Reference 12)

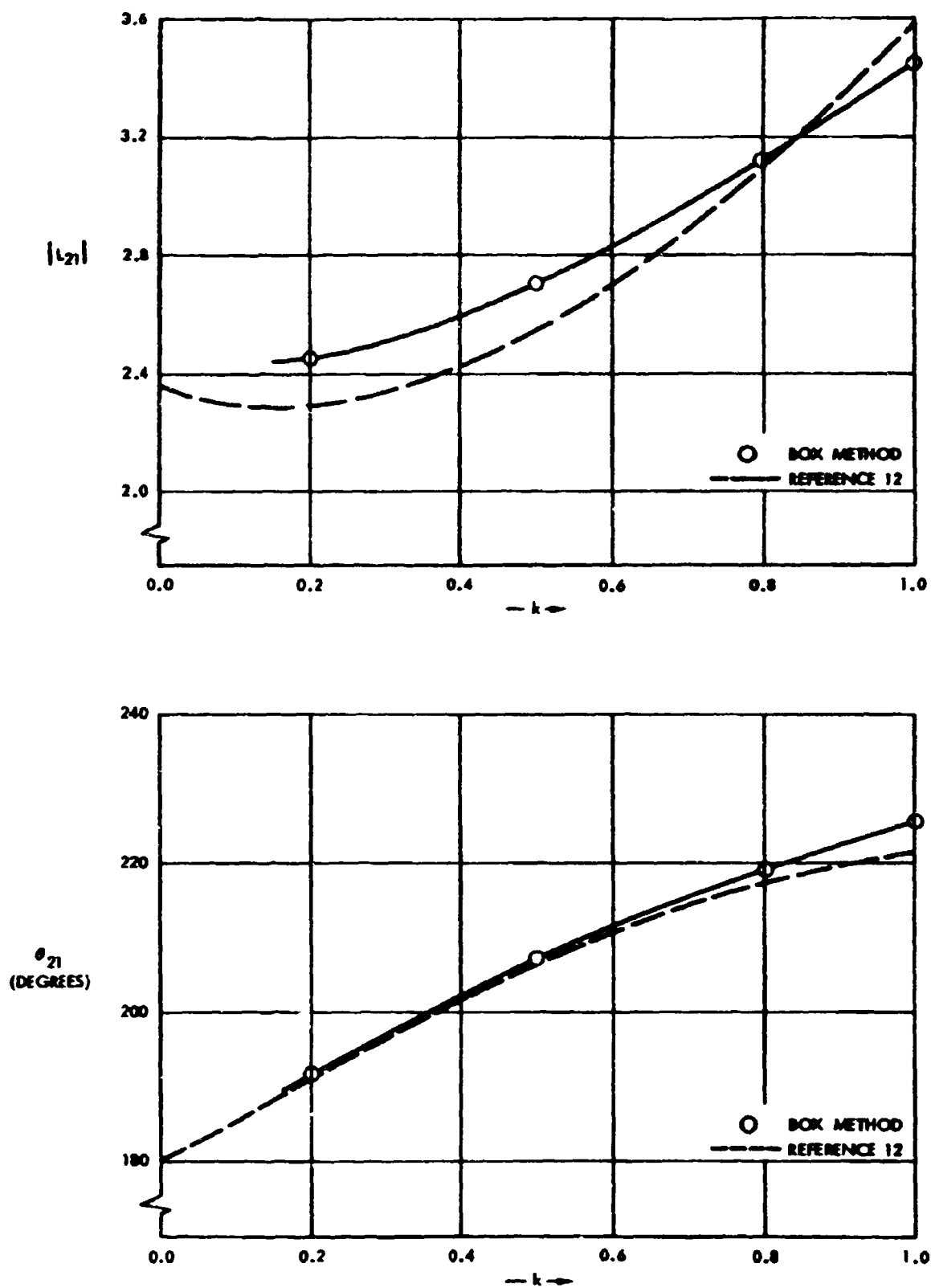


Figure 6. Lift Due to Pitch for an Aspect Ratio 1.5 Delta Wing  
(Compared With Reference 12)



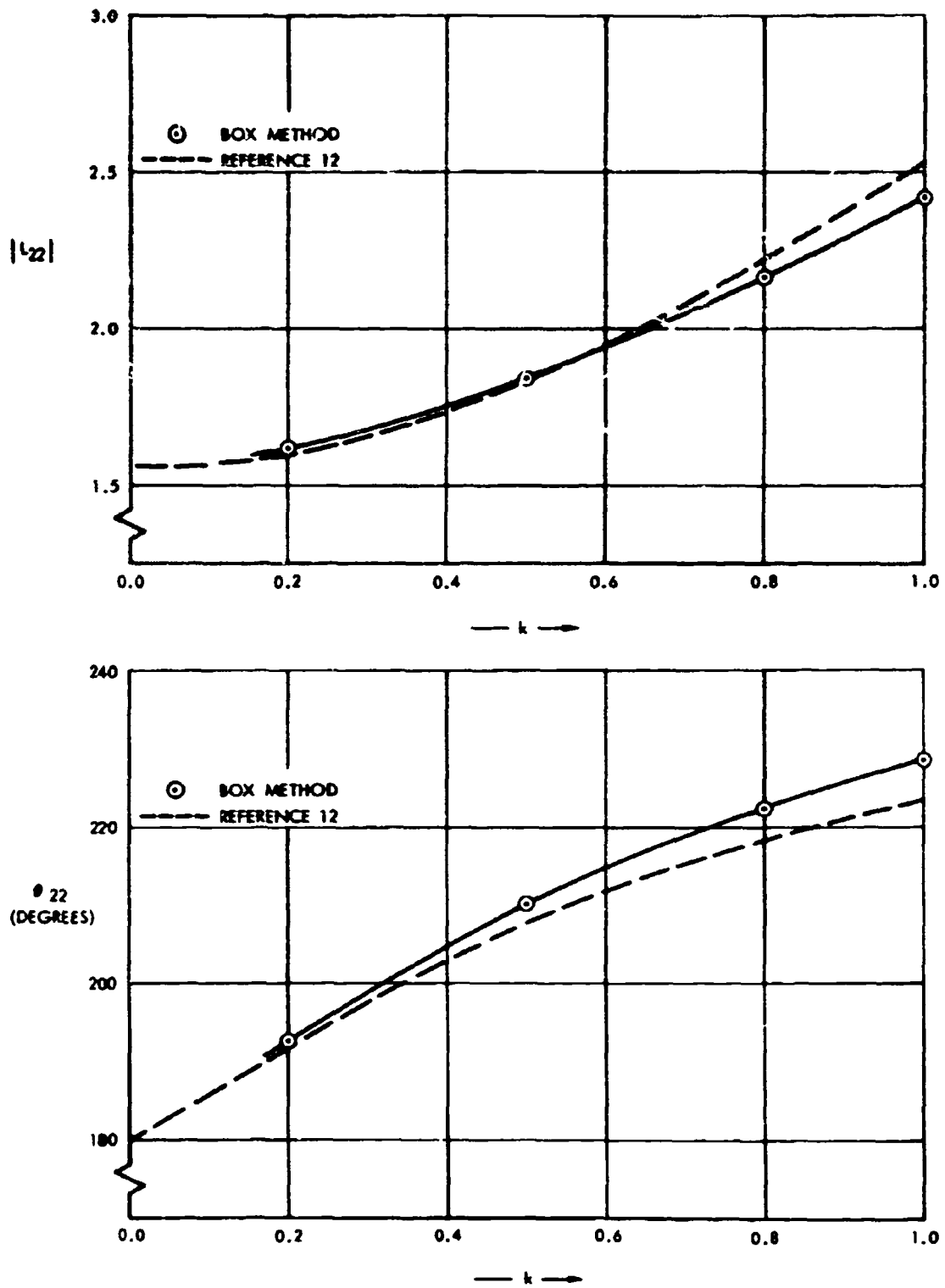


Figure 7. Moment Due to Pitch for an Aspect Ratio 1.5 Delta Wing (Compared With Reference 12)

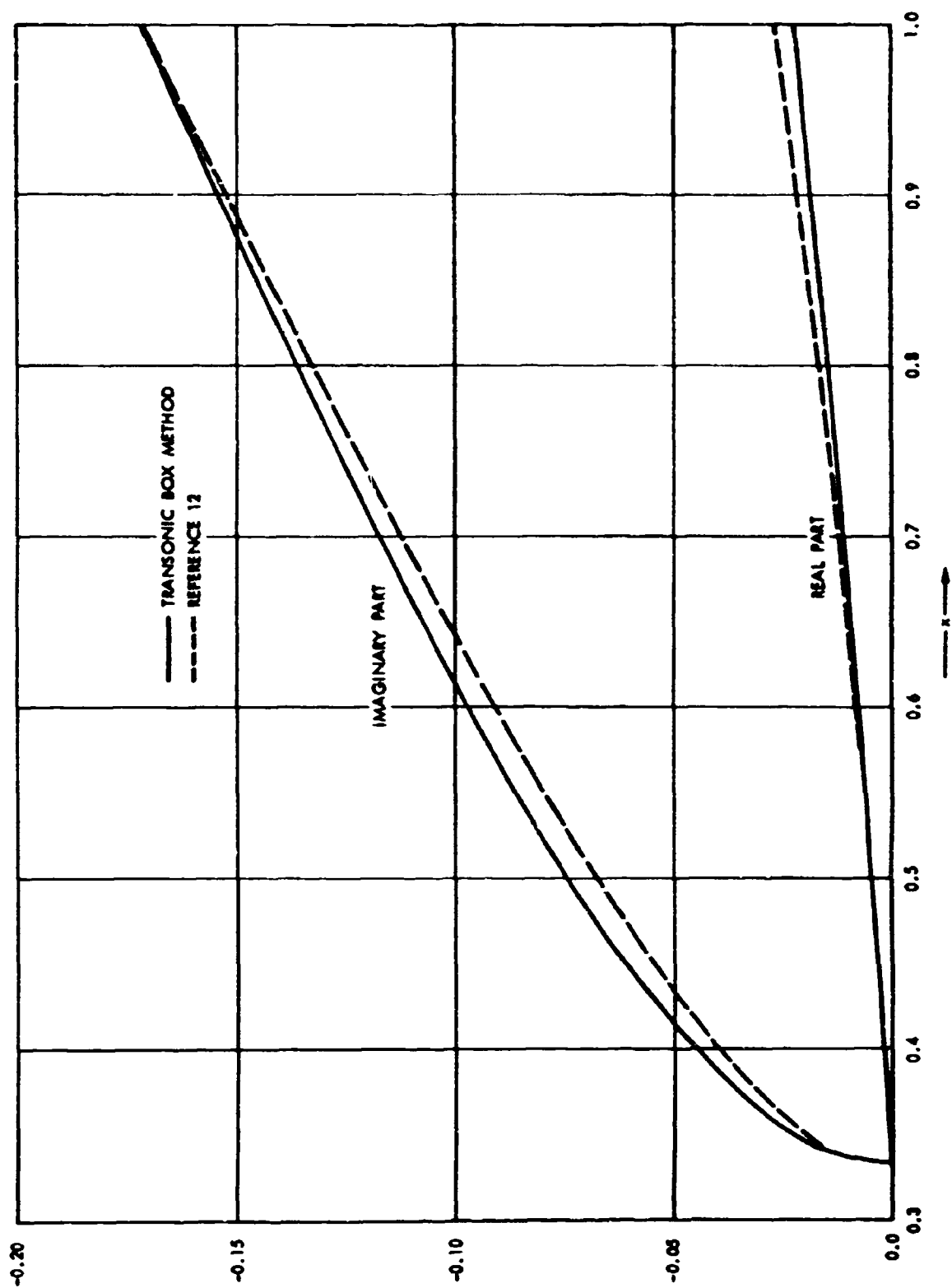


Figure 8. Real and Imaginary Parts of the Unsteady Potential  $\bar{\phi}$  in the  
Plunging Mode for an Aspect Ratio 1.5 Delta Wing at  
 $k = 0.5$ ; Chordwise Distribution for  $y = 0.125 = g \text{ Max}/3$   
(Compared With Reference 12)

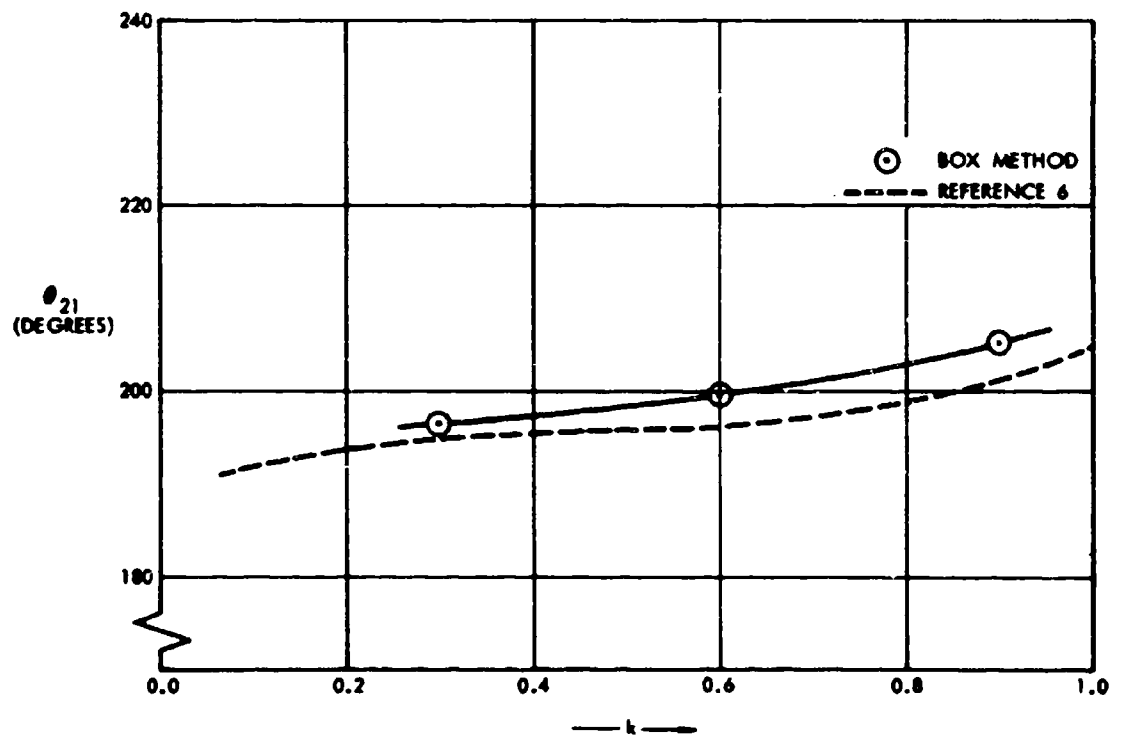
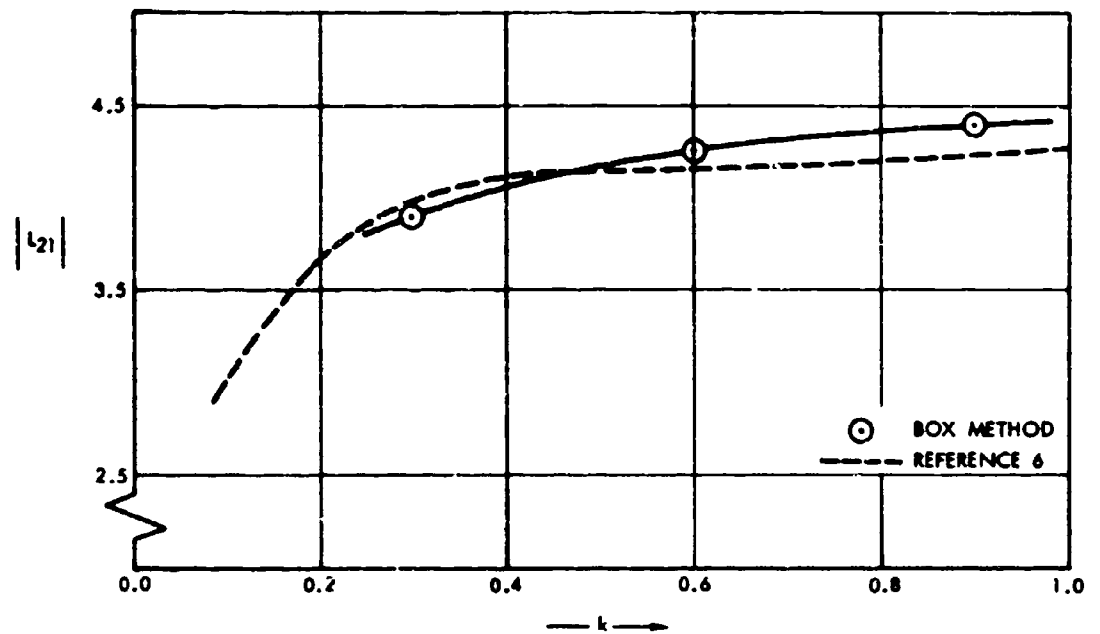


Figure 9. Lift Due to Pitch for an Aspect Ratio 2.0 Rectangular Wing  
(Compared With Reference 6)

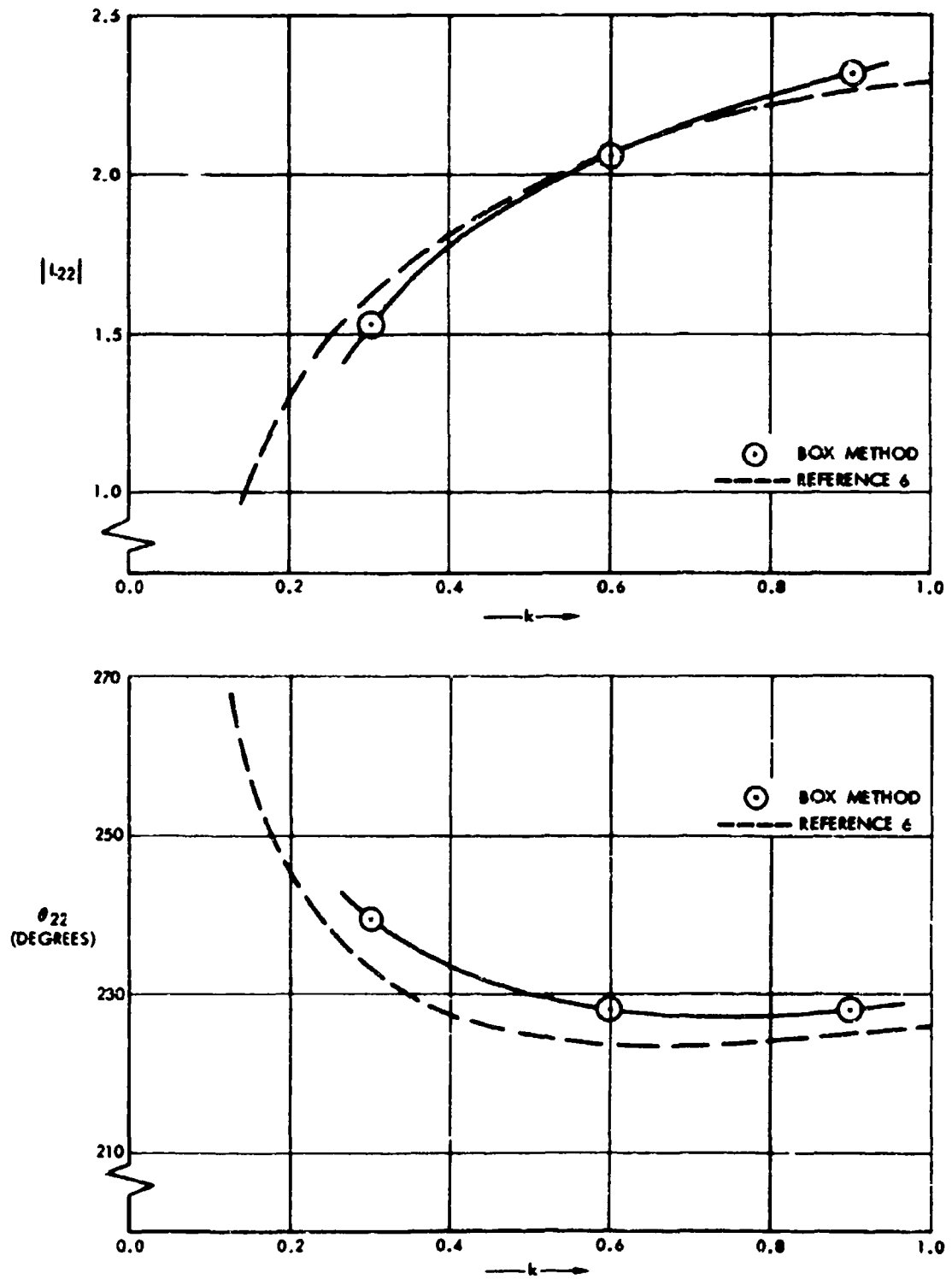


Figure 10. Moment Due to Pitch for an Aspect Ratio 2.0 Rectangular Wing  
(Compared With Reference 6)

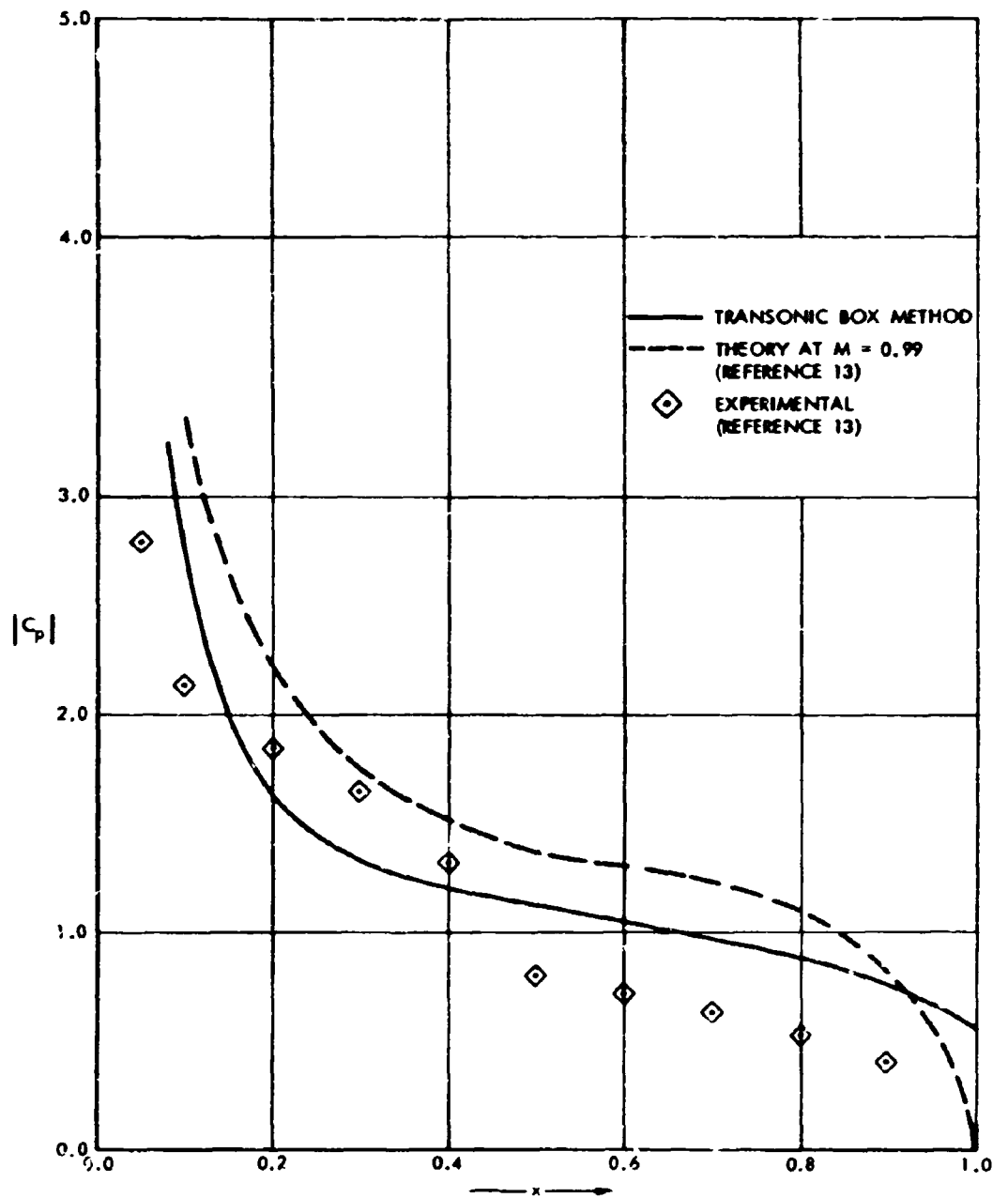


Figure 13. Chordwise Pressure Distribution on an Aspect Ratio 3.0 Rectangular Wing in an Elastic Mode;  $y = y_{\max}/2 = 0.75$  (Compared With Reference 13)

The variation of the experimental values from the computed values is of the type that thickness effects should be expected to cause: the measured pressure is (1) smaller near the leading edge, (2) larger before the point of maximum thickness ( $x = 0.5$ ), and (3) smaller beyond this point. The experimentally determined values of local Mach number along this chord range from 0.84 to 1.35, which indicates that the thickness has a considerable effect on the flow.

The theoretical curve given in Reference 13 was obtained from the subsonic kernel function method, applied at  $M = 0.99$ . This curve is included to show how another theoretical method compares with the experimental values.

#### 4. COMPUTER RUNNING TIME

The results described in this section required about 20 minutes total computer time on the IBM 7094. With nonessential output omitted, this time could have been reduced. All optional output was given, resulting in about 40,000 lines of output.

## 5. CONCLUSIONS

A procedure has been developed for predicting unsteady aerodynamic forces and pressures on an oscillating wing by the use of the transonic box method. The results obtained by this method agree quite well with theoretical values from other methods that are applicable only to special planforms. The box method has the advantage of applicability to a general planform. The only other method of this generality at Mach 1 is the sonic limit of the subsonic kernel function method (see Reference 16) that has not been very successful.

The comparison with experimental values in Figure 11 indicates that the most serious limitation of the method is that thickness is neglected. Thickness may be incorporated into a box program by using modified forms for sources and doublets, depending on the local Mach number (see Reference 14). This was not accomplished under the present program. Other possible extensions of the transonic box method, that would not require much change in the existing computer program, are described in Paragraph 10 of Section 2.

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## APPENDIX I. PROPERTIES OF SOURCE AND DOUBLET DISTRIBUTIONS

### 1. BOUNDARY BEHAVIOR OF A DOUBLET DISTRIBUTION

We wish to evaluate

$$\bar{\varphi}_d(x, y, 0+) = \lim_{z \rightarrow 0+} \iint_{\xi > 0} d\xi d\eta \rho(\xi, \eta) \bar{\varphi}_1(x-\xi, y-\eta, z).$$

The integrand is zero for  $\xi > x$ . If we define  $\rho(\xi, \eta) = 0$  for  $\xi < 0$ , then

$$\bar{\varphi}_d(x, y, 0+) = \lim_{z \rightarrow 0+} \iint_{\xi < x} d\xi d\eta \rho(\xi, \eta) \bar{\varphi}_1(x-\xi, y-\eta, z).$$

Put

$$\xi = x - z^2 u$$

$$\eta = y - zv$$

Then, using Equation (14), we have

$$\begin{aligned} \bar{\varphi}_d(x, y, 0+) &= \lim_{z \rightarrow 0+} \iint_{u>0} du dv \rho(x-z^2 u, y-zv) \cdot z^3 \bar{\varphi}_1(z^2 u, zv, z) \\ &= \lim_{z \rightarrow 0+} \iint_{u>0} du dv \rho(x-z^2 u, y-zv) \cdot \frac{ik}{2\pi} \frac{1}{z^2} e^{-\frac{1}{2} ik \left( z^2 u + \frac{1+v^2}{u} \right)} \end{aligned}$$

Let  $(x, y)$  be a point at which  $\rho$  is continuous. Then the value of  $\rho$  in the integrand approaches  $\rho(x, y)$  as  $z \rightarrow 0$ . Taking the limit under the integral sign,

$$\bar{\varphi}_d(x, y, 0+) = \rho(x, y) \cdot \frac{ik}{2\pi} \iint_{u>0} du dv \cdot \frac{1}{u^2} e^{-\frac{1}{2} ik \frac{1+v^2}{u}}$$

Let  $v = s \sqrt{u}$ . Then

$$\bar{\varphi}_d(x, y, 0+) = \rho(x, y) \cdot \frac{ik}{2\pi} \int_0^{\infty} \frac{du}{u^{3/2}} e^{-\frac{1}{2}ik/u} \int_{-\infty}^{\infty} ds e^{-\frac{1}{2}iks^2}$$

These integrals may be reduced to a standard form by rotating the paths integration in the complex plane into positions in which the exponents are negative, then making the substitutions

$$u = \frac{ik}{2p}$$

$$s = \sqrt{\frac{2}{ik}} q$$

The result is

$$\begin{aligned} \bar{\varphi}_d(x, y, 0+) &= \rho(x, y) \cdot \frac{1}{\pi} \int_0^{\infty} e^{-p} p^{-\frac{1}{2}} dp \int_{-\infty}^{\infty} e^{-q^2} dq \\ &= \rho(x, y) \end{aligned}$$

since both integrals have the value  $\sqrt{\pi}$  (see Reference 12, formulas 507, 512). Consequently, Equation (18) is valid at any point of continuity of  $\rho(x, y)$ .

## 2. BOUNDARY BEHAVIOR OF A SOURCE DISTRIBUTION

To evaluate  $\bar{\varphi}_{sz}(x, y, 0+)$ , note that

$$\begin{aligned} \frac{\partial \bar{\varphi}_s}{\partial z} &= \iint_{\xi > 0} d\xi d\eta \rho(\xi, \eta) \frac{\partial}{\partial z} \bar{\varphi}_0(x-\xi, y-\eta, z) \\ &= \iint_{\xi > 0} d\xi d\eta \rho(\xi, \eta) \bar{\varphi}_1(x-\xi, y-\eta, z) \\ &= \bar{\varphi}_d \end{aligned}$$

Hence, by the result of the preceding section, if  $(x, y)$  is a point at which  $\rho(x, y)$  is continuous,

$$\overline{\varphi}_{sz}(x, y, 0+) = \overline{\varphi}_d(x, y, 0+) = \rho(x, y)$$

which verifies Equation (17).

## APPENDIX II. EXPRESSIONS FOR THE INFLUENCE COEFFICIENTS

Equation (31) may be expressed more conveniently in terms of

$$u = (x_i - \xi)/d$$

$$v = (y_j - \eta)/d$$

$$m = |j - j'|$$

$$n = i - i'$$

$$l = kd$$

( $d$  = box side length). By Equation (14) we have

$$A(n, m) = \frac{ik}{2\pi} \iint \frac{du dv}{u} e^{-\frac{1}{2} i l \left(u + \frac{v^2}{u}\right)} \quad (39)$$

$$\begin{aligned} &|v - m| < \frac{1}{2} \\ &|u - n| < \frac{1}{2} \\ &u \geq 0 \end{aligned}$$

It is assumed that  $l$  is small. If  $l < 0.1$ , the following approximation gives an error of less than 0.1 percent in the value of  $A$ :

$$\begin{aligned} e^{-\frac{1}{2} i l u} &= e^{-\frac{1}{2} i l n} e^{-\frac{1}{2} i l (u - n)} \\ &\approx e^{-\frac{1}{2} i l n} \left[ 1 - \frac{1}{2} i l (u - n) - \frac{1}{8} l^2 (u - n)^2 \right] \end{aligned}$$

This reduces Equation (39) to

$$\begin{aligned}
 A(n, m) = \frac{ik}{2\pi} e^{-\frac{1}{2} i \ell n} & \left\{ \left( 1 + \frac{1}{2} i \ell n - \frac{1}{8} \ell^2 n^2 \right) \iint \frac{du dv}{u^2} e^{-\frac{1}{2} i \ell v^2/u} \right. \\
 & + \left( -\frac{1}{2} i \ell + \frac{1}{4} \ell^2 n \right) \iint \frac{du dv}{u} e^{-\frac{1}{2} i \ell v^2/u} \\
 & \left. - \frac{1}{8} \ell^2 \iint du dv e^{-\frac{1}{2} i \ell v^2/u} \right\}
 \end{aligned} \tag{40}$$

where the limits of integration are the same as in Equation (39).

The following formula expresses these double integrals in terms of single integrals:

$$\begin{aligned}
 \int_{u_1}^{u_2} du \int_{v_1}^{v_2} dv \frac{1}{u^p} e^{-\frac{1}{2} i \ell v^2/u} &= \frac{1}{3-2p} \int_{u_1}^{u_2} \frac{v du}{u^p} e^{-\frac{1}{2} i \ell v^2/u} \Big|_{v=v_1}^{v_2} \\
 &+ \frac{2}{3-2p} \int_{v_1}^{v_2} \frac{dv}{u^{p-1}} e^{-\frac{1}{2} i \ell v^2/u} \Big|_{u=u_1}^{u_2}
 \end{aligned}$$

Equation (40) becomes

$$A(n, m) = \frac{ik}{2\pi} e^{-\frac{1}{2} i \ell n} \cdot \begin{cases} A_n(n + \frac{1}{2}, m) - A_n(n - \frac{1}{2}, m), & n > 0 \\ A_0(\frac{1}{2}, m), & n = 0 \end{cases} \tag{41}$$

where

$$\begin{aligned}
 A_n(u, m) = & v \int_0^u du e^{-\frac{1}{2} i \ell v^2 / u} \left[ -\frac{1}{u^2} \left( 1 + \frac{1}{2} i \ell n - \frac{1}{8} \ell^2 n^2 \right) \right. \\
 & + \frac{1}{u} \left( -\frac{1}{2} i \ell + \frac{1}{4} \ell^2 n \right) - \frac{1}{24} \ell^2 \left. \right] \Big|_{v = m - \frac{1}{2}}^{m + \frac{1}{2}} \\
 & + \left[ -\frac{2}{u} \left( 1 + \frac{1}{2} i \ell n - \frac{1}{8} \ell^2 n^2 \right) + 2 \left( -\frac{1}{2} i \ell + \frac{1}{4} \ell^2 n \right) \right. \\
 & \left. - \frac{1}{12} \ell^2 u \right] \int_{m - \frac{1}{2}}^{m + \frac{1}{2}} dv e^{-\frac{1}{2} i \ell v^2 / u} \\
 = & B_n(u, m) + C_n(u, m)
 \end{aligned} \tag{42}$$

$B_n$  and  $C_n$  denote the contributions of the terms containing the  $u$  - integrals and  $v$  - integrals, respectively.

$B_n(u, m)$  may be expressed in terms of the sine and cosine integrals

$$S(x) = \int_1^{\infty} \frac{\sin xt}{t} dt$$

$$C(x) = \int_1^{\infty} \frac{\cos xt}{t} dt$$

( $C(x)$  and  $S(x)$  are evaluated by the subroutine CIN. ) The resulting formula for  $B_n$  is

$$\begin{aligned}
B_n(u, m) = & \left\{ \left[ \frac{2i}{\ell v} \left( 1 + \frac{1}{2} i \ell n - \frac{1}{8} \ell^2 n^2 \right) - \frac{1}{24} \ell^2 uv \right] e^{-\frac{1}{2} i \ell v^2 / u} \right. \\
& + \left[ v \left( -\frac{1}{2} i \ell + \frac{1}{4} \ell^2 n \right) + \frac{1}{48} i \ell^3 v^3 \right] \left[ C\left(\frac{\ell v^2}{2u}\right) \right. \\
& \left. \left. - i S\left(\frac{\ell v^2}{2u}\right) \right] \right\} \Big|_{v = m - \frac{1}{2}}^{m + \frac{1}{2}}
\end{aligned} \quad (43)$$

To evaluate  $C_n(u, m)$ , put  $v = s + m$ . Expanding part of the exponential gives the approximation

$$\begin{aligned}
\int_{m-\frac{1}{2}}^{m+\frac{1}{2}} dv e^{-\frac{1}{2} i \ell v^2 / u} &= \int_{-\frac{1}{2}}^{\frac{1}{2}} ds e^{-\frac{1}{2} i \ell (m^2 + 2ms + s^2) / u} \\
&\cong e^{-\frac{1}{2} i \ell m^2 / u} \int_{-\frac{1}{2}}^{\frac{1}{2}} ds e^{-i \ell ms / u} \left( 1 - \frac{1}{2} i \ell s^2 / u \right)
\end{aligned}$$

and performing the integration gives

$$\begin{aligned}
\int_{m-\frac{1}{2}}^{m+\frac{1}{2}} dv e^{-\frac{1}{2} i \ell v^2 / u} &= e^{-\frac{1}{2} i \ell m^2 / u} \left\{ \frac{\sin(\ell m / 2u)}{\ell m / 2u} \left( 1 - \frac{1}{8} \frac{i \ell}{u} \right) \right. \\
&\quad \left. + \frac{i u}{\ell m^2} \left[ \frac{\sin(\ell m / 2u)}{\ell m / 2u} - \cos(\ell m / 2u) \right] \right\}, \quad m \neq 0
\end{aligned}$$

For small values of  $\ell m / 2u$ , the trigonometric functions of this argument are expanded in power series. To sufficient accuracy,

$$\int_{m-\frac{1}{2}}^{m+\frac{1}{2}} dv e^{-\frac{1}{2} i \ell v^2 / u} = e^{-\frac{1}{2} i \ell m^2 / u} \left[ 1 - \frac{1}{6} \left( \frac{\ell m}{2u} \right)^2 - \frac{1}{24} \frac{i \ell}{u} \right], \quad \ell m / 2u < 0.2$$



It may be verified that this is valid for  $m = 0$ .

Combining these expressions,

$$C_n(u, m) = e^{-\frac{1}{2} i \ell m^2 / u} \left[ -\frac{2}{u} \left( 1 + \frac{1}{2} i \ell n - \frac{1}{8} \ell^2 n^2 \right) + 2 \left( -\frac{1}{2} i \ell + \frac{1}{4} \ell^2 n \right) - \frac{1}{12} \ell^2 u \right] \cdot \begin{cases} \left( 1 - \frac{1}{2} \frac{i \ell}{u} \right) \frac{\sin(\ell m / 2u)}{\ell m / 2u} + \frac{i u}{\ell m^2} \left[ \frac{\sin(\ell m / 2u)}{\ell m / 2u} - \cos(\ell m / 2u) \right], & \ell m / 2u > 0.2 \\ 1 - \frac{1}{24} \frac{i \ell}{u} - \frac{1}{6} \left( \frac{\ell m}{2u} \right)^2, & \ell m / 2u < 0.2 \end{cases} \quad (44)$$

The subroutine PØT2 evaluates the influence coefficients according to Equations (41), (42), (43), and (44).

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### APPENDIX III. COMPUTER PROGRAM LISTINGS

# MAIN PROGRAM

```

DIMENSION A(2,100,16),W(2,50,50),DA(700),PS(2,5,50),DF(5,50)
DIMENSION ML(50),AXY(9,9),AY(9),XEDG(5),YEDG(5),COE(5,5)
DIMENSION G(10)
COMMON A,W,DA,PS,DF,ML,AXY,AY,XEDG,YEDG,COE,M,L,NS,D,DI,CK,IEDG
COMMON AREA,DH
DATA 2/1H /

INITIALIZATION OF DATA ARRAY

DO 1 I=1,700
  1 DA(I)=0.0
DO 11 I=1,24
  11 DA(I)=2
DO 13 I=104,700,4
  13 DA(I)=1.0

16 CALL DATRD(DA)
CK=DA(23)*DA(24)/DA(25)*6.28318531
DI=DA(27)
L=DI
D=1.0/DI
DH=0.5*D
WRITE
  0 (6, 49)(DA(I),I=1,12)
36 IF (L) 83,83,37
37 IF (50-L) 83,38,38
38 WRITE
  0 (6, 41)(L,CK)
41 FORMAT(1H010X,12,23H BOXES ALONG ROOT CHORD/1H010X,19HREDUCED FREQUENCY =F6.3)
WRITE
  0 (6, 42)(DA(23)
42 FORMAT(1H011X,10HFREQUENCY=1PE11.3)
  IF (DA(26)) 33,35,33
33 DO 34 I=13,22
34 DA(I)=2
GO TO 27

```

```

00000420
00000430
00000440
00000450
00000460
00000470
00000480
00000490
00000500
00000502
00000504
00000510
00000520
00000530
00000540
00000550
00000560
00000570
00000580
00000590
00000600
00000610
00000620
00000630
00000640
00000650
00000660
00000670
00000680
00000690
00000700
00000710
00000720
00000730
00000740
00000750
00000760

```

```

35 CALL SHAPE
   CALL FORCI
27 LIM=ML(L)
   IF (LIM-50) 22,22,101
22 LIM2=2*ML
   LPOT= PINO(L,15)
   CALL POT2(100,LIM2,LPOT,CK,0)
   M=0
   K=DA(28)
   GO TO 4

C
C
C
C
PRELIMINARY CALCULATIONS ARE FINISHED.
THE NEXT SECTION IS GONE THROUGH FOR EACH MODE.

3 IF (DA(26)) 4,28,4
28 CALL DATRD(DA)
4 K=K-1
  M=M+1
  WRITE
    0 (6, 48)M
48 FORMAT (1M115X,8MMODE NO.13)
  WRITE
    0 (6, 49)(DA(I),I=13,22)
49 FORMAT (1M010X,12A6)
  IF (DA(26)) 29,7,29
7 CALL DRED
9 G(M)=DA(39)
29 IF (G(M)) 74,72,74
72 CALL WVAL
  IF (ML) 24,26,24
24 LIM2=2*ML
  DO 25 J=1,LIM2
25 W(J,1,1)=W(J,1,1)+2.0/3.14159265
  LEADING EDGE CORRECTION
26 CONTINUE
  CALL BOXP
74 IF (K) 6,6,5

```

C  
C  
C  
C

C

00000770  
00000780  
00000790  
00000800  
00000810  
00000820  
00000830  
00000840  
00000850  
00000860  
00000870  
00000880  
00000890  
00000900  
00000910  
00000920  
00000930  
00000940  
00000950  
00000960  
00000970  
00000980  
00000990  
00001000  
00001010  
00001020  
00001030  
00001040  
00001050  
00001060  
00001070  
00001080  
00001090  
00001100  
00001110  
00001120  
00001130

```

C
C
C
5 IF (M-10) 3,6,6
    FINAL SECTION OF PROGRAM - COMPUTATION OF GENERALIZED FORCES
6 WRITE (6,46)
46 FORMAT (1H110X,18HGENERALIZED FORCES/1H05X,5HMODES/4X,11H0SC. DEF00001200
    1L.8X,9HREAL PART10X,9HIMAG PART10X,10HABS. VALUE6X,11HPHASE ANGLE)00001210
    AC=8.0/AREA
    DO 12 M1=1,M
    IF (G(M1)) 12,14,12
14 DO 18 M2=1,M
    S1=0.0
    S2=0.0
    N1=5*(M1-1)
    N2=5*(M2-1)
    DO 8 J=1,5
    J1=J+N1
    J2=J+N2
    DO 8 I=1,5
    S1=S1+PS(1,I,J1)*DF(1,J2)
    S2=S2+PS(2,I,J1)*DF(1,J2)
    S1=AC*S1
    S2=AC*S2
    S3= SQRT(S1**2+S2**2)
    S4= ATAND(S2,S1)
18 WRITE
    0 (6, 47)M1,M2,S1,S2,S3,S4
47 FORMAT (1H02I6,1P3E19.5,OP1F16.4)
12 CONTINUE
    WRITE ( 6,43)
43 FORMAT(1H1)
    GO TO 16
C
C
C
    ERROR EXITS
83 IPR=27
84 WRITE
00001140
00001150
00001160
00001170
00001180
0000001200
0000001210
00001220
00001230
00001240
00001250
00001260
00001270
00001280
00001290
00001300
00001310
00001320
00001330
00001340
00001350
00001360
00001370
00001380
00001390
00001400
00001410
00001420
00001430
00001440
00001460
00001470
00001480
00001490
00001500
00001510
00001520

```

```

0 (6, 45)IPR
45 FORMAT(1M010X, 8H8AD DATA14)
GO TO 102
101 WRITE ( 6, 56)
56 FORMAT(1M010X, 42HLATERAL LIMIT ON NUMBER OF BOXES EXCEEDED.)
102 STOP
END

```

```

00001530
00001540
00001550
00001560
00001580
00001600
00001610

```

# SUBROUTINE SHAPE

## SUBROUTINE SHAPE

```

DIMENSION A(2,100,16),W(2,50,50),DA(700),PS(2,5,50),DF(5,50)
DIMENSION ML(50),AXY(9,9),AY(9),XEDG(5),YEDG(5),COE(5,5)
COMMON A,W,DA,PS,DF,ML,AXY,AY,XEDG,YEDG,COE,M,L,NS,D,DI,CK,IEDG
COMMON AREA,DH
IEDG=0
NS=DA(29)
IF (NS) 81,81,1
1 IF (NS-3) 2,2,81
2 NSP=NS+1
IF (DA(24)) 82,82,3
3 DO 4 I=1,NSP
XEDG(I)=DA(2+I+27)/DA(24)
4 YEDG(I)=DA(2+I+28)/DA(24)
XEDG=0.0
XEDG(NS+2)=1.0
YEDG(NS+2)=YEDG(NS+1)
Y1=DA(30)
IF (Y1) 83,20,20
20 K=0
X=DH
AREA=0.0
DO 15 I=1,1
7 IF (X-XEDG(K+1)) 8,8,9
8 ML(I)=0.5+DI*(YEDG(K)+F*(X-XEDG(K)))/G)
GO TO 15
9 K=K+1
G=XEDG(K+1)-XEDG(K)
F=YEDG(K+1)-YEDG(K)
IF (G) 84,12,12
12 IF (F) 85,13,13
13 AREA=AREA+G*(YEDG(K+1)+YEDG(K))
GO TO 7
15 X=X+D
IF (XEDG(NS+1)-1.0) 17,16,84
16 IF (YEDG(NS+1)-YEDG(NS)) 17,17,18
17 IEDG=1

```

00001640  
00001650  
00001660  
00001670  
00001680  
00001690  
00001700  
00001710  
00001720  
00001730  
00001740  
00001750  
00001760  
00001770  
00001780  
00001790  
00001800  
00001810  
00001820  
00001830  
00001840  
00001850  
00001860  
00001870  
00001880  
00001890  
00001900  
00001910  
00001920  
00001930  
00001940  
00001950  
00001960  
00001970  
00001980  
00001990  
00002000



# SUBROUTINE SHAPE

```

18 RETURN
81 IPR=29
   GO TO 86
82 IPR=24
   GO TO 86
83 IPR=30
   GO TO 86
84 K= HINC(K,NS)
   IPR=2*K+29
   GO TO 86
85 IPR=2*K+30
86 WRITE
   0 (6, 41)IPR
41 FORMAT(1H010X,8HBAU DATA13)
   STOP
   END

```

```

00002010
00002020
00002030
00002040
00002050
00002060
00002070
00002080
00002090
00002100
00002110
00002120
00002130
00002140
00002160
00002170

```

```

SUBROUTINE FORCI
CONTROL SUBROUTINE FOR CALCULATION OF INTEGRALS USED
TO FIND GENERALIZED FORCES

SUBROUTINE FORCI
DIMENSION A(2,100,16),W(2,50,50),DA(700),PS(2,5,50),DF(5,50)
DIMENSION ML(50),AXY(9,9),AY(9),XEDG(5),YEDG(5),COE(5,5)
COMMON A,W,DA,PS,DF,ML,AXY,AY,XEDG,YEDG,COE,M,L,NS,D,DI,CK,IEDG
COMMON AREA,DH
COMMON YMAX2,N
YMAX2=YEDG(NS+1)**2
DO 3 J=1,9
DO 2 I=1,9
2 AXY(I,J)=0.0
3 AY(J)=0.0
4 IF (DA(30)) 4,5,4
4 N=0
CALL SECT
SECT DOES THE CALCULATIONS FOR EACH SECTION OF THE PLANFORM

5 DO 7 I=1,NS
IF (YEDG(I)-YEDG(I+1)) 6,7,7
6 N=I
CALL SECT
7 CONTINUE
RETURN
END

```

```

00002200
00002210
00002220
00002230
00002240
00002250
00002260
00002270
00002280
00002290
00002300
00002310
00002320
00002330
00002340
00002350
00002360
00002370
00002380
00002390
00002400
00002410
00002420
00002430
00002440
00002450
00002460
00002470

```

[illegible]

SUBROUTINE SECT

H(4)=H(3)  
H(5)=H(2)  
H(6)=H(1)

GAUSSIAN POINTS AND WEIGHTS FOR THE INTERVAL (0,1)

X2=XEDG(N+1)  
X1=XEDG(N)  
Y2=YEDG(N+1)  
Y1=YEDG(N)  
IF (N) 4,2,4

2 X1=0.0

Y1=0.0

4 DY=Y2-Y1

DO 19 J=1,6

V=U(J)

G=H(J)\*DY

IF (Y2\*\*2-YMAX2) 7,6,7

6 G=2.0\*V\*G

V=V\*V

7 Y=Y2-V\*DY

X0=X2+V\*(X1-X2)

XOP=1.0-X0

G=G\* SQRT(XOP)

YQ=Y\*Y

IF (IEDG) 8,9,8

8 G=G\* SQRT(1.0-YQ/YMAX2)

9 E=2.0\*XOP\*G

IF (DA(30)) 10,10,11

10 G=G\* SQRT(1.0+X0)

11 DO 17 I=1,6

U2=U(I)\*\*2

X=X0+XOP\*U2

F=E\*H(I)\*U2

IF (DA(30)) 12,12,13

12 F=F\* SQRT(X+X0)

13 YP=1.0

00002870  
00002880  
00002890  
00002900  
00002910  
00002920  
00002930  
00002940  
00002950  
00002960  
00002970  
00002980  
00002990  
00003000  
00003010  
00003020  
00003030  
00003040  
00003050  
00003060  
00003070  
00003080  
00003090  
00003100  
00003110  
00003120  
00003130  
00003140  
00003150  
00003160  
00003170  
00003180  
00003190  
00003200  
00003210  
00003220  
00003230

# SUBROUTINE SECT

```

00 16 M=1,9
XP=YP
00 15 L=1,9
AXY(L,M)=AXY(L,M)+XP*F
15 XP=X*XP
16 YP=YQ*YP
17 CONTINUE
YP=L.O
00 18 M=1,9
AY(M)=AY(M)+YP*G
18 YP=YQ*YP
19 CONTINUE
RETURN
END

```

```

00003240
00003250
00003260
00003270
00003280
00003290
00003300
00003310
00003320
00003330
00003340
00003350
00003360
00003370

```

```

SUBROUTINE DRED
SUBROUTINE DRED
DIMENSION A(2,100,16),W(2,50,50),DA(700),PS(2,5,50),DF(5,50)
DIMENSION ML(50),AXY(9,9),AY(9),XEDG(5),YEDG(5),COE(5,5)
COMMON A,W,DA,PS,DF,ML,AXY,AY,XEDG,YEDG,COE,M,L,NS,D,DI,CK,IEOG
COMMON AREA,OH
DIMENSION I(8(5),B(25),C(25,25),G(25)
COMMON C,B,G,IXB
NP=DA(98)
IF (NP) 83,24,30
C
C
C
C
      A POLYNOMIAL FOR THE DEFLECTION IS FITTED TO VALUES
      OF DEFLECTION AT GIVEN POINTS.
30 NX=DA(99)
   NY=DA(100)
   IF (NX) 81,81,31
   IF (NY) 82,82,32
31 IF (NY) 82,82,32
32 IF (150-NP) 83,38,38
38 IF (NP-NX) 81,34,34
34 IF (NP-NY) 82,35,35
35 MX= MINO(NX,5)
   MY= MINO(NY,5)
   IY=MY
   IXY=MX+MY
   DO 2 J=1,5
   DO 1 I=1,5
      COE(I,J)=0.0
      1 IXB(J)=MX
      NC=MX+MY
      3 DO 5 I=1,NC
      DO 4 J=1,NC
      C(I,J)=C.0
      5 B(I)=0.0
      KP=100
      DO 11 IP=1,NP
      X=DA(KP+1)/DA(24)
      Y2=(DA(KP+2)/DA(24))*2
00003400
00003410
00003420
00003430
00003440
00003450
00003460
00003470
00003480
00003490
00003500
00003510
00003520
00003530
00003540
00003550
00003560
00003570
00003580
00003590
00003600
00003610
00003620
00003630
00003660
00003662
00003664
00003670
00003680
00003690
00003700
00003710
00003720
00003730
00003740
00003750
00003760

```

# SUBROUTINE DRED

```

DEF=DA(KP+3)
WT=DA(KP+4)
IF (WT) 84,84,6
6 YP=1.0
K=1
DO 8 J=1,IY
XYP=YP
JX=IXB(J)
DO 7 I=1,JX
G(K)=XYP
XYP=X+XYP
7 K=K+1
8 YP=Y2+YP
DO 10 I=1,NC
DO 9 J=1,NC
9 C(I,J)=C(I,J)+G(I)*G(J)*WT
10 B(I)=B(I)+G(I)*DEF*WT
11 KP=KP+4
K=MSIMER(25,NC,1,C,B)
IF (K-1) 22,22,15
15 DO 16 I=1,IY
IP=IY+1-I
IF (IXB(IP)+IP-IXY) 16,17,17
16 CONTINUE
IXY=IXY-1
GO TO 15
17 IXB(IP)=IXB(IP)-1
IF (IXB(IP)) 18,18,19
18 IY=IP-1
19 NC=0
DO 20 I=1,IY
20 NC=NC+IXB(I)
GO TO 3
22 K=1
DO 23 J=1,IY
JX=IXB(J)
DO 23 I=1,JX

```

00003770  
00003780  
00003790  
00003800  
00003810  
00003820  
00003830  
00003840  
00003850  
00003860  
00003870  
00003880  
00003890  
00003900  
00003910  
00003920  
00003930  
00003940  
00004030  
00004040  
00004050  
00004060  
00004070  
00004080  
00004090  
00004100  
00004110  
00004120  
00004130  
00004140  
00004150  
00004160  
00004170  
00004180  
00004190  
00004200  
00004210

SUBROUTINE DRED

```

C0E(I,J)=0(K)
23 K=K+1
IF (DA(07)) 61,66,61
61 WRITE (6,41)
DO 64 J=1,IV
JX=IX0(J)
DO 63 I=1,JX
WRITE
0 (6, 42)I,J,C0E(I,J)
63 CONTINUE
64 CONTINUE
66 CONTINUE
GO TO 28
24 YP=1.0
K=1
DO 27 J=1,5
XYP=YP
DO 26 I=1,5
C0E(I,J)=XYP*DA(K+45)
K=K+1
26 XYP=XYP*DA(24)
27 YP=YP*DA(24)**2
28 K=25*(M-1)
DO 29 I=1,25
K=K+1
29 DF(K,1)=C0E(1,1)
RETURN
81 IPR=99
GO TO 85
82 IPR=100
GO TO 85
83 IPR=98
GO TO 85
84 IPR=K+4
85 WRITE
0 (6, 45)IPR
STOP

```

```

00004220
00004230
00004240
00004250
00004270
00004280
00004290
00004300
00004310
00004320
00004330
00004340
00004350
00004360
00004370
00004380
00004390
00004400
00004410
00004420
00004430
00004440
00004450
00004460
00004470
00004480
00004490
00004500
00004510
00004520
00004530
00004540
00004550
00004560
00004570
00004580
00004600

```



SUBROUTINE ORED

```

41 FORMAT(1H010X,56HCOMPUTED DEFLECTION = SUM OF DEF(N,M)*X*(N-1)*Y*00004610
1*(2M-2)/1H010X,54H(IN DIMENSIONLESS COORDINATES - DISTANCE/CHORD L 00004620
2ENGTH)/1H09X,1HN7X,1HM16X,8HDEF(N,M)
42 FORMAT(3X,218,1PE25.5)
45 FORMAT(1H010X,8HBAC DATA14)
END
00004630
00004640
00004650
00004660

```

SUBROUTINE WVAL

SUBROUTINE WVAL

EVALUATION OF THE UPWASH ARRAY

DIMENSION A(2,100,16),W(2,50,50),DA(700),PS(2,5,50),DF(5,50)  
DIMENSION ML(50),AXY(9,9),AY(9),XEDG(5),YEDG(5),COE(5,5)  
COMMON A,W,DA,PS,DF,ML,AXY,AY,XEDG,YEDG,COE,M,L,NS,D,DI,CK,IEDG  
COMMON AREA,DM

DIMENSION G(5,5),H(5,5)

COMMON G,H

J1=5\*(M-1)

DO 3 J=1,5

J1=J1+1

CI=1.0

DO 2 I=1,5

G(I,J)=CI\*DF(I+1,J1)

H(I,J)=CK\*G(I,J1)

2 CI=CI+1.0

3 G(5,J)=0.0

G AND H ARE THE COEFFICIENTS OF THE REAL AND IMAGINARY PARTS  
OF THE UPWASH

X=DM

DO 10 I=1,L

JL=ML(I)

IF (JL) 5,10,5

5 Y=DM

DO 9 J=1,JL

Y2=Y\*Y

W(1,J,I)=0.0

W(2,J,I)=0.0

YP=1.0

DO 8 J1=1,5

XYP=Y\*Y

DO 7 I1=1,5

W(1,J,I)=W(1,J,I)+G(I1,J1)\*XYP

00004690

00004700

00004710

00004720

00004730

00004740

00004750

00004760

00004770

00004780

00004790

00004800

00004810

00004820

00004830

00004840

00004850

00004860

00004870

00004880

00004890

00004900

00004910

00004920

00004930

00004940

00004950

00004960

00004970

00004980

00004990

00005000

00005010

00005020

00005030

00005040

00005050

SUBROUTINE HVAL

W(2,J,I)=W(2,J,I)+H(I1,J1)\*XYP  
7 XYP=X\*XYP  
8 YP=Y2\*YP  
9 Y=Y+D  
10 X=X+D  
RETURN  
END

00005060  
00005070  
00005080  
00005090  
00005100  
00005110  
00005120

# SUBROUTINE BOXP

```

SUBROUTINE BOXP
  DIMENSION A(2,100,16),S(2,50,50),DA(700),PS(2,5,50),DF(5,50)
  DIMENSION ML(50),AXY(9,9),AY(9),XEDG(5),YEDG(5),COE(5,5)
  DIMENSION EDG(50),PR(2,50),PSI(2,50),G(20),XO(50),IXB(4)
  DIMENSION C(20,20),B(20,2)
  COMMON A,S,DA,PS,DF,ML,AXY,AY,XEDG,YEDG,COE,M,L,NS,O,DI,CK,IEDG
  COMMON AREA,OM
  COMMON B,C,EDG,PR,PSI,G,XO,IXB

  IF (IDA(09)) 71,75,71
71 WRITE (6,47)
47 FORMAT(1H11OX,43HTIME UPWASH ARRAY (REAL AND IMAGINARY PARTS))
D0 74 I=1,L
  JL=ML(I)
  IF (JL) 73,74,73
73 WRITE
    0 (6, 42)I
  IF (I-1) 77,77,79
77 D0 78 J=1,JL
  S1=S(1,J,I)+1.57079633
  S2=S(2,J,I)+1.57079633
78 WRITE
    0 (6, 142)S1,S2
  GO TO 74
79 WRITE
    0 (6, 142)(S(1,J,I),S(2,J,I),J=1,JL)
142 FORMAT(1H 1P2E24.5)
74 CONTINUE
  THESE ARE THE UPWASHES
C
C
75 CONTINUE
  CALL BOXPO
  BOXPO COMPUTES THE POTENTIAL VALUES IN EACH BOX.
  THEY ARE STORED IN THE ARRAY S.
C
C
  IF (IDA(09)) 91,95,91
91 WRITE (6,143)

```

# SUBROUTINE B0XP

143 FORMAT(1H11OX.46HTHE POTENTIAL ARRAY (REAL AND IMAGINARY PARTS))

DO 94 I=1,L

JL=ML(I)

IF (JL) 93,94,93

93 WRITE

0 (6, 42)I

WRITE

0 (6, 142)(S(1,J,I),S(2,J,I),J=1,JL)

94 CONTINUE

THESE ARE THE POTENTIALS

95 CONTINUE

FIT OF A SERIES TO THE POTENTIAL VALUES

JL=ML(L)

DY=D/YEDG(NS+1)

Y=0.5\*DY

DO 201 J=1,JL

IF (IEDG) 202,203,202

202 EDG(J)= SQRT(1.0-Y\*Y)

Y=Y+DY

GO TO 201

203 EDG(J)=1.0

201 CONTINUE

N=0.5\*DI\*YEDG

IF (N) 4,4,2

2 DO 3 I=1,N

3 X0(I)=0.0

4 X1=0.0

N1=N

Y1=YEDG

DO 8 K=1,NS

X2=XEDG(K+1)

Y2=YEDG(K+1)

N=N1\*Y2+0.5

IF (N1-N) 5,7,7

5 N1=N1+1

SUBROUTINE BCXP

```

      DO 6 I=1,N
      Y=0*( FLGAT(I)-0.5)
      6 XD(I)=X1+(X2-X)*(Y-Y1)/(Y2-Y1)
      7 X1=X2
      Y1=Y2
      8 N1=N
      AS=0.0
      DO 119 I=1,L
      JL=ML(I)
      IF (JL) 119,119,217
      217 DO 118 J=1,JL
      AS=AMAX1(AS, ABS(S(1,J,I)), ABS(S(2,J,I)))
      118 CONTINUE
      119 CONTINUE
      IX=5
      IY=4
      IXY=9
      13 DO 14 I=1,IY
      14 IXB(I)=IX
      16 NC=0
      DO 17 I=1,IY
      17 NC=NC+IXB(I)
      DO 19 I=1,NC
      DO 18 J=1,NC
      C(I,J)=0.0
      B(I,1)=0.0
      19 B(I,2)=0.0
      X1=0.5*0
      DO 25 I=1,L
      JL=ML(I)
      Y=0.5*0
      IF (JL) 25,25,20
      20 DO 24 J=1,JL
      XR=X1-X0(J)
      IF (YEDG) 601,601,602
      601 XR=XR*(X1+X0(J))
      602 XR= SQRT(XR)

```

# SUBROUTINE BOXP

```

YP=1.0
K=1
DO 23 N1=1,IY
XP=XR*YP
JX=IXB(N1)
DO 22 N=1,JX
G(K)=XP*EDG(J)
XP=X1*XP
22 K=K+1
23 YP=Y*Y*YP
Y=Y+D
DO 24 N1=1,NC
DO 124 N=1,NC
124 C(N1,N)=C(N1,N)+G(N1)*G(N)
DO 24 N=1,2
24 B(N1,N)=B(N1,N)+G(N1)*S(N,J,1)
25 X1=X1+D
K=MSIMER(20,NC,2,C,B)
IF (K-1) 30,30,29
C
C
C
C
IF XSIMQ FAILS, THE FOLLOWING SECTION REDUCES THE NUMBER OF TERMS
IN THE SERIES.
29 DO 61 I=1,IY
IP=IY+1-I
IF (IXB(IP)+IP-IXY) 61,62,62
61 CONTINUE
IXY=IXY-1
GO TO 29
62 IXB(IP)=IXB(IP)-1
IF (IXB(IP)) 63,63,16
63 IY=IY-1
GO TO 16
C
30 K=i
AC=0.0
DO 133 I=1,IY

```

00006740  
00006750  
00006760  
00006764  
00006766  
00006770  
00006780  
00006785  
00006790  
00006800  
00006810  
00006820  
00006830  
00006840  
00006850  
00006860  
00006870  
00006880  
00006890  
00006900  
00006910  
00006920  
00006940  
00006950  
00006970  
00006980  
00007000  
00007010  
00007030  
00007040  
00007060  
00007080  
00007090  
00007100  
00007110  
00007120  
00007140

```

C      JX=IX8(I)
C      CJ=1.0
C      DO 33 J=1,JX
C      C(K,1)=B(K,1)
C      C(K,2)=B(K,2)
C      AC=AMAX1(AC, ABS(C(K,1)), ABS(C(K,2)))
C      IF (J.EQ.1) GO TO 33
C      B(K-1,1)=CJ*C(K,1)
C      B(K-1,2)=CJ*C(K,2)
C      CJ=CJ+1.0
C      33 K=K+1
C      B(K-1,1)=0.0
C      133 B(K-1,2)=0.0
C      IF (AC-50.0*AS) 117,117,29
C      GO TO 29 IF COEFFICIENTS ARE TOO LARGE
C      117 IF (DA(90)) 110,113,110
C      PRINTOUT OF COEFFICIENTS OF POTENTIAL
C      110 WRITE (6,44)
C      IF (YEDG) 603,603,605
C      603 WRITE (6,45)
C      IF (IEDG) 604,607,604
C      604 WRITE (6,48)
C      GO TO 607
C      605 WRITE (6,145)
C      IF (IEDG) 606,607,606
C      606 WRITE (6,148)
C      607 WRITE (6,46)
C      K=1
C      DO 111 I=1,IY
C      JL=IX8(I)
C      DO 111 J=1,JL
C      WRITE (6,49) J,I,C(K,1),C(K,2)
C      111 K=K+1

```



# SUBROUTINE BOXP

```

      WRITE ( 6,149)
113 IF (DA(91)) 114,116,114
C
C      PRINTOUT OF VALUES OF POTENTIAL AND PRESSURE
C
114 WRITE ( 6,41)
41 FORMAT(1M010X,9HPOTENTIAL45X,8MPRESSURE)
X1=0.5*D
DO 39 I=1,L
JL=ML(I)
Y=0.5*D
IF (JL) 39,39,34
34 DO 38 J=1,JL
XR=X1-X0(J)
XQ=0.5
IF (YEDG) 608,608,609
608 XQ=X1
XR=XR*(X1+X0(J))
609 XQ=XQ/XR
XR= SQRT(XR)
DO 37 N=1,2
PSI(N,J)=0.0
PR(N,J)=0.0
K=1
YP=EDG(J)
DO 37 N1=1,IY
XPI=XR*YP
JX=IXB(N1)
DO 36 M=1,JX
PSI(N,J)=PSI(N,J)+C(K,N)*XPI
PR(N,J)=PR(N,J)+B(K,N)*XPI
XPI=X1*XPI
36 K=K+1
37 YP=Y+Y*YP
PR(1,J)=2.0*(PR(1,J)+PSI(1,J)*XQ-PSI(2,J)*CK)
PR(2,J)=2.0*(PR(2,J)+PSI(2,J)*XQ+PSI(1,J)*CK)
38 Y=Y+D

```

```

0000715G
00007170
00007180
00007190
00007200
00007210
00007230
00007240
00007250
00007260
00007270
00007280
00007290
00007300
00007310
00007320
00007330
00007340
00007350
00007360
00007370
00007380
00007390
00007400
00007410
00007420
00007430
00007440
00007450
00007460
00007470
00007480
00007490
00007500
00007510
00007520
00007530

```

SUBROUTINE ROXP

```

      WRITE
      0 (6, 42)I
      42 FORMAT(1H010X,12,6HTH ROW)
      WRITE
      0 (6, 43)(PSI(1,J),PSI(2,J),PK(1,J),PR(2,J),J=1,JL)
      43 FORMAT(1H 1P2E24.5,5X,2E24.5)
      39 X1=X1+D
C
      116 JMO=5*(M-1)
      CJ=0.0
      DO 172 I=1,5
      DO 72 J=1,5
      J0=J+JMO
      PS (1,I,J0)=0.0
      PS (2,I,J0)=0.0
      K=1
      DO 72 K2=1,IY
      J2=J+K2-1
      JX=IXB(K2)
      DO 72 K1=1,JX
      J1=K1+I-1
      X1=AY(J2)-CJ*AXY(J1-1,J2)
      Y=AXY(J1,J2)*CK
      PS (1,I,J0)=PS (1,I,J0)+C(K,1)*X1-C(K,2)*Y
      PS (2,I,J0)=PS (2,I,J0)+C(K,1)*Y +C(K,2)*X1
      72 K=K+1
      172 CJ=CJ+1.0
      RETURN
C
      44 FORMAT(1H-10X,53HPOTENTIAL = SUM OF PG(M,N)*X***(M-1)*Y***(2N-2)*SQR00007830
      1TF(X)
      45 FORMAT(1H+63X,10H**2-X0**2))
      145 FORMAT(1H+63X,4H-X0))
      48 FORMAT(1H+73X,21H*SQRFF(1-(Y/YMAX)**2))
      148 FORMAT(1H+67X,21H*SQRFF(1-(Y/YMAX)**2))
      46 FURMAT(1H015X,52HWHERE X = X0(Y) IS THE EQUATION OF THE LEADING ED00007890
      1GE./1H020X,21HCOEFFICIENTS PG(M,N)/1H07X,1HM7X,1HN14X,9HREAL PART00007900

```

SUBROUTINE BEXP

216X,10HIMAG. PART)  
49 FORMAT(1H 218,1P2E25.5)  
149 FORMAT(1H1)  
END

00007910  
00007920  
00007930  
00007940

```

SUBROUTINE BOXPO
C
C SOLUTION OF SIMULTANEOUS EQUATIONS FOR THE POTENTIAL
COMMON A(2,100,16),S(2,50,50),DA(700),PS(2,5,50),DF(5,50),ML(50)
COMMON AX(9,9),AY(9),XEDG(5),YEDG(5),CDE(5,5),M,L,NS,D,OI,CK,IEDG
COMMON AREA,DH,E(2,50,50)
I1=0
DO 9 I=1,L
KO=MAXO(1,I-14)
JL=ML(I)
IF (JL.EQ.O) GO TO 9
IF (I1.EQ.O) GO TO 6
C SUBTRACTION OF CONTRIBUTIONS OF PRECEDING ROWS TO UPWASH
DO 5 J=1,JL
DO 5 K=KO,I1
KL=ML(K)
K1=I+1-K
IF (KL.EQ.O) GO TO 5
DO 4 N=1,KL
N1=N+J
N2=IABS(N-J)+1
A1=A(1,N1,K1)+A(1,N2,K1)
A2=A(2,N1,K1)+A(2,N2,K1)
S(1,J,I)=S(1,J,I)-A1*S(1,N,K)+A2*S(2,N,K)
4 S(2,J,I)=S(2,J,I)-A2*S(1,N,K)-A1*S(2,N,K)
5 CONTINUE
C SETTING UP MATRIX FOR SIMULTANEOUS EQUATIONS
DO 8 J=1,JL
DO 8 K=1,J
N1=J+K
N2=IABS(J-K)+1
E(1,J,K)=A(1,N1,1)+A(1,N2,1)
E(2,J,K)=A(2,N1,1)+A(2,N2,1)
E(1,K,J)=E(1,J,K)
8 E(2,K,J)=E(2,J,K)
C SOLUTION OF EQUATIONS
K=MSIMEC(50,JL,1,E,S(1,1,I))
IF (K.NE.1) GO TO 12
9 I1=I1+1

```

SUBROUTINE BOXPO

RETURN

12 WRITE ( 6,41)

41 FORMAT(1H010X,59HSOLUTION OF SIMULTANEOUS EQUATIONS FOR THE POTENT

IAL FAILED)

STOP

END

10390  
10400  
10410  
10420  
10430  
10440

C  
C  
C  
C  
C  
C  
C  
C  
C  
C

# SUBROUTINE POT2

THE VELOCITY FIELD OF A UNIFORM DOUBLE DISTRIBUTION  
OVER A BOX IS COMPUTED AT ALL POINTS AT WHICH IT WILL BE  
NEEDED AND STORED IN THE ARRAY A IN COMMON

MO,NO CONTROL THE NUMBER OF VALUES COMPUTED

M2 IS THE RANGE OF THE SECOND SUBSCRIPT IN THE ARRAY,  
DIMENSIONED A(2,M2,N2), BUT TREATED HERE AS AN ARRAY  
WITH TWO SUBSCRIPTS

SUBROUTINE POT2(M2,MO,NO,CK,U)  
DIMENSION A(2,1)

COMMON A

M=MO

N=NO

OK=CK\*0

OK2=OK\*\*2

M1=M-1

OK8=OK2/8.0

OK4=2.0\*OK8

OK12=OK2/12.0

CM=0.5

DM=OK\*0.5

DM=0.5\*DM

DD=2.0\*CK

DDM=DD

D1=0.25\*OK2

B5=OK2/24.0

DO 3 I=1,M

B1=0.0

B4=2.0/DM

B2=B5/B4-DH

B3=-0.5\*B5

D3=DM\*B4+B5

D4=OK8\*B4

DD4=2.0\*DD4

00007970  
00007980  
00007990  
00008000  
00008010  
00008020  
00008030  
00008040  
00008050  
00008060  
00008070  
00008080  
00008090  
00008100  
00008110  
00008120  
00008130  
00008140  
00008150  
00008160  
00008170  
00008180  
00008190  
00008200  
00008210  
00008220  
00008230  
00008240  
00008250  
00008260  
00008270  
00008280  
00008290  
00008300  
00008310  
00008320  
00008330

# SUBROUTINE POT2

```

CN=1.0
K=I
C3=0.0
C4=0.0
C7=0.0
C8=0.0
DO 2 J=1,N
  AI=DM/CN
  C1=CM* COS(A1)
  C2=-CM* SIN(A1)
  C5=CM*CIN(A1,C6)
  C6=-CM*C6
  C9=C1-C3
  C10=C2-C4
  C11=C5-C7
  C12=C6-C8
  A(1,K)=B3=C9-B4=C10-B5=C3-B1=C11-B2=C12
  A(2,K)=B4=C9+B3=C10-B5=C4+B2=C11-B1=C12
23 C3=C1
  C4=C2
  C7=C5
  C8=C6
  B1=B1-D1
  B3=B3-D3
  B4=B4-D4
  D4=D4+D04
  CN=CN+2.0
2 K=K+M2
  CM=CM+1.0
  DM=DM+DDM
3 DDM=DDM+DD
  D5 5 L=1,2
  K1=1
DO 5 J=1,N
DO 4 I=1,M1
  K=K1+M-I
4 A(L,K)=A(L,K)-A(L,K-1)

```

```

00008340
00008350
00008360
00008370
00008380
00008390
00008400
00008410
00008420
00008430
00008440
00008450
00008460
00008470
00008480
00008490
00008500
00008510
00008520
00008530
00008540
00008550
00008560
00008570
00008580
00008590
00008600
00008610
00008620
00008630
00008640
00008650
00008660
00008670
00008680
00008690
00008700

```

SUBROUTINE POT2

```

      A(L,K1)=2.0*A(L,K1)
5  K1=K1+M2
      CM=0.0
      DM=0.0
      DDM=DK
      DO 12 I=1,M
      C7=0.0
      C8=0.0
      C9=0.0
      C10=0.0
      P1=0.0
      P2=0.0
      CN=1.0
      B6=0.5*OK12
      K=I
      DO 10 J=1,N
      A1=CM/CN
      A2=DM/CA
      IF (A1-0.2) 7,7,8
7  B1=2.0-A1*2/3.0
      B2=-DK/(6.0*CN)
      GO TO 9
8  B3= SIN(A1)/A1
      B1=2.0*B3
      B2=(B3- COS(A1))/A2-DH/CN*B3
9  B3= COS(A2)/CN
      B4= SIN(A2)/CN
      C3=B1*B3+B2*B4
      C4=B2*B3-B1*B4
      B5=DH*CN
      C1=B5*C4-2.0*C3
      C2=-2.0*C4-B5*C3
      C5=C1-C7
      C6=C2-C8
      P3=P2-B6*CN
      P4=P3+2.0*OK12*(CN-1.0)
      A(1,K)=A(1,K)+C5-P1*C6+P3*C3-P4*C9

```

00008710  
 00008720  
 00008730  
 00008740  
 00008750  
 00008760  
 00008770  
 00008780  
 00008790  
 00008800  
 00008810  
 00008820  
 00008830  
 00008840  
 00008850  
 00008860  
 00008870  
 00008880  
 00008890  
 00008900  
 00008910  
 00008920  
 00008930  
 00008940  
 00008950  
 00008960  
 00008970  
 00008980  
 00008990  
 00009000  
 00009010  
 00009020  
 00009030  
 00009040  
 00009050  
 00009060  
 00009070



SUBROUTINE PGT2

```

A(2,K)=A(2,K)+C6+P1+C5+P3+C4-P4+C10
P1=P1+DH
P2=P2+CN+DK4
CN=CN+2.0
C7=C1
C8=C2
C9=C3
C10=C4
B6=B6+DK12
10 K=K+M2
CM=CM+DK
DM=DM+DDH
12 ODM=ODM+DD
D3=CK/(2.0+3.14159265)
M1=M2-M
K=1
A1=0.0
DO 14 J=1,N
C1=D3* SIN(A1)
C2=-D3* COS(A1)
DO 13 I=1,M
DF =A(1,K)+C1+A(2,K)*C2
A(2,K)=A(2,K)+C1-A(1,K)*C2
A(1,K)=DF
13 K=K+1
K=K+M1
14 A1=A1+DH
RETURN
END

```

```

00009080
00009090
00009100
00009110
00009120
00009130
00009140
00009150
00009160
00009170
00009180
00009190
00009200
00009210
00009220
00009230
00009240
00009250
00009260
00009270
00009280
00009290
00009300
00009310
00009320
00009330
00009340
00009350
00009360

```

```

CIN(X,S)
SINE AND COSINE INTEGRAL SUBROUTINE
IF CALLED BY THE STATEMENT C=CIN(X,S)
C AND S ARE THE INTEGRALS OVER T FROM 1 TO INFINITY OF
COS(XT)/T AND SIN(XT)/T
FUNCTION CIN(X1,S)
SG=1.0
X=X1
IF (X) 1,2,2
1 SG=-SG
X=-X
2 X2=X*X
IF (X-1.0) 3,3,4
FOR ABS(X) LESS THAN 1 A SERIES EXPANSION IS USED
3 V=((X2/98.0-0.6)*.05*X2+1.0)*X2/18.0-1.0)*X+1.57079633
U=((X2/45.0-1.0)*X2/24.0+1.0)*X2/4.0-.577215665-ALOG(X)
GO TO 5
FOR ABS(X) GREATER THAN 1 APPROXIMATIONS OF HASTINGS ARE USED
4 P=((X2+19.394119)*X2+47.411538)*X2+8.493336)/(((X2+21.361055)
1 *X2+70.376456)*X2+30.038227)*X)
Q=((X2+21.383724)*X2+49.719775)*X2+5.089504)/(((X2+27.177958)
1 *X2+119.918932)*X2+76.707876)*X2)
CO=COS (X)
SI=SIN (X)
U=Q*CO-P*SI
V=P*CO+Q*SI
5 S=V*SG
CIN=U
RETURN
END

```

DATRO

```

C      CARD-READ SUBROUTINE 'DATRO(DATA(I))'
SUBROUTINE DATRO(DATA)
  DIMENSION ORBU(14),DATA(1)
  DATA ATEST/5HALPHA/,DTEST/1H /,ETEST/1H-/
  1 READ (5,2) EMIN,ALP,IND,DRBU(1),I=1,12)
  2 FORMAT(A1,A5,16,12A6)
  IF (ALP.EQ.ATEST) GO TO 9

C
  WRITE (99,2) EMIN,ALP,IND,DRBU(1),I=1,12)
  CARD IS WRITTEN IN INTERNAL BUFFER
  REWIND 99
  IF (ALP.NE.DTEST) GO TO 8

C
  NUMERIC CARD

  READ (99,990) DRBU(1),I=1,5)
  REWIND 99
  DO 5 I=1,5
    IF(DRBU(1))4,6,4
    DATA(IND)=DRBU(1)
    IND=IND+1
  5 GO TO 11

C
  TEST FOR BLANK FIELD
  IF(SIGN(1.0,DRBU(1)))5,5,4

C
  ALPHA CARD

  DO 10 I=1,10
    DATA(IND)=DRBU(1)
    IND=IND+1
  10 IF (EMIN.NE.ETEST) GO TO 1

C
  RETURN IF COLUMN 1 CONTAINS A MINUS SIGN

  13 RETURN
C

```

00021410  
00021420  
00021430  
00021435  
00021450  
00021455  
00021460  
00021465  
00021470  
00021475  
00021480  
00021490  
00021500  
00021510  
00021520  
00021530  
00021540  
00021550  
00021560  
00021570  
00021580  
00021590  
00021600  
00021610  
00021620  
00021630  
00021640  
00021650  
00021770  
00021780  
00021790  
00021800  
00021805  
00021810  
00021815  
00021820  
00021825

DATRD

C  
C

BAD CARD

8 READ (99,992) DRBU  
WRITE ( 6,993) DRBU  
WRITE ( 6,991)  
REWIND 99  
STOP

990 FORMAT(12X,5E12.0)  
991 FORMAT(38H BAD DATA ON THIS CARD. JOB TERMINATED )  
992 FORMAT(14A6)  
993 FORMAT(12HOCARD IMAGE=14A6)  
END

00021830  
00021835  
00021840  
00021845  
00021850  
00021855  
00021860  
00021870  
00021880  
00021890  
00021900  
00021910

# SIMULTANEOUS EQUATION SUBROUTINE

```

* 18MAP SIMR      K=MSIMER(N,L,LB,A,B)
* SOLVES THE SYSTEM OF EQUATIONS  A*X=B.
* TO USE, SET K=MSIMER(N,L,LB,A,B)
* WHERE N IS THE NUMBER OF ROWS FOR WHICH A IS
* DIMENSIONED, AND L IS THE NUMBER OF EQUATIONS.
* LB IS THE NUMBER OF COLUMNS IN B.
* K=1 DENOTES SUCCESSFUL SOLUTION
* K=2 FOR A SINGULAR OR ILL-CONDITIONED MATRIX
* K=3 IF IMPROPER DATA IS GIVEN.
* TO AVOID THIS SIGNAL, L MUST BE POSITIVE AND AT MOST 100,
* N MUST NOT BE LESS THAN L, A MUST NOT INCLUDE A ROW
* OF ZEROS.
* A IS DESTROYED.
* IF K=1, THE SOLUTION IS RETURNED IN B
*
* ENTRY
*
* MSIMER SAVE  1,2,3,4,5,6,7
*
* PROLOGUE
*
* CLA= 3,4
* PAX  0,1
* TXL  EL,1,0
* TXM  EL,1,100
* PCO  0,1
* STD  A4-1
* STD  A6-1
* STD  2A6-1
* STD  A16
* STD  A16+1
* STD  2A21
* STD  A12+1
* STD  2A26
* STD  2A26+1
* STD  2A37
* STD  2A37+1

```

```

00050000
00050020
00050030
00050040
00050050
00050060
00050070
00050080
00050090
00050100
00050110
00050120
00050130
00050140
00050150
00050160
00050170
00050180
00050190
00050200
00050210
00050220
00050230
00050240
00050250
00050260
00050270
00050280
00050290
00050300
00050310
00050320
00050330
00050340
00050350
00050360
00050370

```

# SIMULTANEOUS EQUATION SUBROUTINE

STD	A12+2	00050380
STD	A20	00050390
STD	A22	00050400
STD	A25+1	00050410
STD	A25+2	00050420
STD	A33	00050430
STD	A36+1	00050440
STD	A36+2	00050450
SXD	A52,1	00050460
TXI	+1,1,1	00050470
SCD	A32,1	00050480
TXI	+1,1,-2	00050490
SCD	A22+1,1	00050500
CLA*	5,4	00050510
PAX	0,7	00050520
SXA	1A6-1,7	00050530
SXA	A14,7	00050540
SXA	1A21-1,7	00050550
SXA	A26,7	00050560
SXA	A37,7	00050570
CLA*	4,4	00050580
PAX	0,1	00050590
TXL	E1,1,0	00050600
TXH	E1,1,**	00050610
SXA	A2,1	00050620
SXA	A5-1,1	00050630
SXD	A9,1	00050640
SXD	A12,1	00050650
TXI	+1,1,-1	00050660
SXD	A7,1	00050670
SXD	A36,1	00050680
SXD	A18,1	00050690
SXD	A38-1,1	00050700
SXD	A21-1,1	00050710
SXD	A23,1	00050720
SXD	A25,1	00050730
SXD	A31,1	00050740

A52

# SIMULTANEOUS EQUATION SUBROUTINE

00050750  
00050760  
00050770  
00050780  
00050790  
00050800  
00050810  
00050820  
00050830  
00050840  
00050850  
00050860  
00050870  
00050880  
00050890  
00050900  
00050910  
00050920  
00050930  
00050940  
00050950  
00050960  
00050970  
00050980  
00050990  
00051000  
00051010  
00051020  
00051030  
00051040  
00051050  
00051060  
00051070  
00051080  
00051090  
00051100  
00051110

SCD A51,1  
SCD A51+1,1  
TXI +1,1,1  
CLA 6,6  
PAC 0,3  
CLA 7,4  
PAC 0,5  
TXI +1,3, -L+1  
TXI +1,5, -L+1

A51  
.  
.  
.  
A2

NORMALIZATION OF ROWS

AXT L,2  
PXA 0,3  
PAX 0,6  
PAX 0,7  
PXA 0,0  
LDQ A,6  
LRS 0  
TLQ +2  
XCA  
-NX  
TXI A4+2,1  
TZE A3,6, -N  
STG E1  
CLA T  
FDP =1,0  
STQ T  
AXT T  
RMP L,2  
STG A,7  
LDQ A,7  
TNX T  
TXI A6+2,1  
PXA A5,7, -N  
PAX 0,5  
AXT 0,6  
LB,4

A3  
A4  
A5  
A6

# SIMULTANEOUS EQUATION SUBROUTINE

1A6	FMP	B,6	00051120
	STG	B,6	00051130
	LQ	T	00051140
	TNX	2A6,4,1	00051150
2A6	TXI	1A6,6,-N	00051160
	TNX	A7,1,1	00051170
	TXI	+1,3,1	00051180
	TXI	A2,5,1	00051190
A7	TXH	B7,2,L -1	00051200
.			00051210
.			00051220
.			00051230
	PXA	0,2	00051240
	PAX	0,1	00051250
	PXA	0,3	00051260
	PAX	0,6	00051270
	PXA	0,0	00051280
A8	LQ	A,6	00051290
	LRS	0	00051300
	TLQ	+3	00051310
	XCA		00051320
	SXA	A10,1	00051330
	TXI	+1,1,1	00051340
	TXH	+2,1,L	00051350
A9	TXI	A8,6,-1	00051360
	LQ	TOL	00051370
	TLQ	+2	00051380
	TRA	E3	00051390
A10	AXT	+1	00051400
	SXD	+1,2	00051410
	TNX	A17,1,++	00051420
.			00051430
.			00051440
.			00051450
	PXA	0,3	00051460
	PAX	0,6	00051470
	PAX	0,7	00051480

SEARCH FOR MAXIMUM PIVOT IN COLUMN

ROW INTERCHANGE



# SIMULTANEOUS EQUATION SUBROUTINE

SCD	++1,1	00051490
TXI	++1,7,++	00051500
PXA	0,2	00051510
PAX	0,4	00051520
CLA	A,6	00051530
LDQ	A,7	00051540
STQ	A,7	00051550
STQ	A,6	00051560
TXI	++1,4,1	00051570
TXH	A13,4,L	00051580
TXI	++1,6, -N	00051590
TXI	A11,7, -N	00051600
PXA	0,5	00051610
PAX	0,6	00051620
PAX	0,7	00051630
AXT	L8,4	00051640
SCD	++1,1	00051650
TXI	++1,6,++	00051660
CLA	B,7	00051670
LDG	B,6	00051680
STQ	B,6	00051690
STQ	B,7	00051700
TXH	A17,4,1	00051710
TXI	++1,6, -N	00051720
TXI	A15,7, -N	00051730
		00051740
		00051750
		00051760
		00051770
		00051780
		00051790
		00051800
		00051810
		00051820
		00051830
		00051840
		00051850

	DIVISION OF ROW BY PIVOT	
A17	CLA	=1,0
	FDP	A,3
	STQ	AM
A18	TXH	A21,2,L -1
	PXA	0,2
	PAX	0,4
	PXA	0,3
	PAX	0,6
A20	TXI	++1,6, -N



# SIMULTANEOUS EQUATION SUBROUTINE

A25	TXI	++1,4,1	00052230
	TXH	A27,4,L -1	00052240
	TXI	++1,3, -N	00052250
	TXI	A24,7, -N	00052260
A27	AXT	++,3	00052270
A26	AXT	LR,4	00052280
	AXT	++,7	00052290
	TXI	++1,7,1	00052300
	SXA	--2,7	00052310
1A26	LDQ	A,6	00052320
	FMP	B,5	00052330
	CHS		00052340
	FAD	B,7	00052350
	STG	B,7	00052360
	TNX	3A26,4,1	00052370
2A26	TXI	++1,5, -N	00052380
	TXI	1A26,7, -N	00052390
3A26	AXT	++,5	00052400
	TNX	A29,1,1	00052410
A28	AXT	++,7	00052420
	PXA	O,2	00052430
	PAX	O,4	00052440
	TXI	++1,3,1	00052450
	TXI	A23,6,1	00052460
A29	AXT	++,3	00052470
A31	TXH	A43,2,L -1	00052480
	PXA	O,2	00052490
	PAX	O,1	00052500
	PAX	O,4	00052510
	PXA	O,3	00052520
	PAX	O,6	00052530
	PAX	O,7	00052540
A32	TXI	++1,3, -N-1	00052550
	TXI	++1,6,-1	00052560
A33	TXI	++1,7, -N	00052570
	SXA	A37+1,5	00052580
	SXA	A41,5	00052590

# SIMULTANEOUS EQUATION SUBROUTINE

A34	SXA	A40,3	00052600
A35	SXA	A39,7	00052610
	SXA	A38,3	00052620
	LDQ	A,6	00052630
	FMP	A,7	00052640
	CHS		00052650
	FAD	A,3	00052660
	STO	A,3	00052670
A36	TXI	+1,4,1	00052680
	TXH	A37,4,L -1	00052690
	TXI	+1,3,-N	00052700
	TXI	A35,7,-N	00052710
A37	AXT	L8,4	00052720
	AXT	*,7	00052730
	TXI	+1,7,-1	00052740
	SXA	-2,7	00052750
1A37	LDQ	A,6	00052760
	FMP	B,5	00052770
	CHS		00052780
	FAD	B,7	00052790
	STO	B,7	00052800
2A37	TNX	A41,4,1	00052810
	TXI	+1,5,-N	00052820
	TXI	1A37,7,-N	00052830
A41	AXT	*,5	00052840
	TXI	+1,1,1	00052850
A38	TXH	A40,1,L -1	00052860
A39	AXT	*,3	00052870
	AXT	*,7	00052880
	PXA	0,2	00052890
	PAX	0,4	00052900
	TXI	+1,3,-1	00052910
	TXI	A34,6,-1	00052920
A40	AXT	*,3	00052930
	TXI	+1,5,-1	00052940
	TXI	A7,2,1	00052950
A43	CLA	=1	00052960

00052970  
00052980  
00052990  
00053000  
00053010  
00053020  
00053030  
00053040  
00053050  
00053060  
00053070  
00053080  
00053090  
00053100  
00053110  
00053120  
00053130  
00053140  
00053150  
00053160  
00053170  
00053180  
00053190  
00053200  
00053210  
00053220  
00053230

95

# COMPLEX SIMULTANEOUS EQUATION SUBROUTINE

```

SIMAP SIMC
*      K=MSIMEC(N,L,LB,A,B)
*      SOLVES THE SYSTEM OF COMPLEX EQUATIONS A*X=B.
*      TO USE, SET K=MSIMEC(N,L,LB,A,B)
*      WHERE N IS THE NUMBER OF ROWS FOR WHICH A IS
*      DIMENSIONED, AND L IS THE NUMBER OF EQUATIONS.
*      LB IS THE NUMBER OF COLUMNS IN B.
*      K=1 DENOTES SUCCESSFUL SOLUTION
*      K=2 FOR A SINGULAR OR ILL-CONDITIONED MATRIX
*      K=3 IF IMPROPER DATA IS GIVEN.
*      TO AVOID THIS SIGNAL, L MUST BE POSITIVE AND AT MOST 100,
*      N MUST NOT BE LESS THAN L, & MUST NOT INCLUDE A ROW
*      OF ZEROS.
*      A IS DESTROYED.
*      IF K=1, THE SOLUTION IS RETURNED IN B
*
*      ENTRY
*
*      PSIMEC SAVE 1,2,3,4,5,6,7
*
*      PROLOGUE
*
CLA* 3,4
ALS 1
PAX 0,1
TXL E1,1,1
TXH E1,1,200
PCD 0,1
STD A4-1
STD A5-1
STD 2A6-1
STD A16
STD A16+1
STD 2A21
STD A12+1
STD 2A26
STD 2A26+1
STD 2A37

```

# COMPLEX SIMULTANECUS EQUATION SUBROUTINE

STD	2A37+1	00053620
STD	A12+2	00053630
STD	A20	00053640
STD	A22	00053650
STD	A25+1	00053660
STD	A25+2	00053670
STD	A33	00053680
STD	A36+1	00053690
STD	A36+2	00053700
SXD	A52+1	00053710
TXI	+1,1,2	00053720
SCD	A32+1	00053730
TXI	+1,1,-4	00053740
SCD	A22+1,1	00053750
CLA*	5,4	00053760
PAX	0,7	00053770
SXA	1A6-1,7	00053780
SXA	A14,7	00053790
SXA	1A21-1,7	00053800
SXA	A26,7	00053810
SXA	A37,7	00053820
CLA*	4,4	00053830
ALS	1	00053840
PAX	0,1	00053850
TXL	E1,1,1	00053860
TXH	E1,1,00	00053870
SXA	A2,1	00053880
SXA	A5-1,1	00053890
SXD	A9,1	00053900
SXD	A12,1	00053910
TXI	+1,1,-2	00053920
SXD	A7,1	00053930
SXD	A36+1	00053940
SXD	A18+1	00053950
SXD	A38-1,1	00053960
SXD	A21-1,1	00053970
SXD	A23,1	00053980

A52

# COMPLEX SIMULTANEOUS EQUATION SUBROUTINE

	SXD	A25,1	00053990
	SXD	A31,1	00054000
	SCD	A51,1	00054010
	SCD	A51+1,1	00054020
	TXI	*+1,1,2	00054030
	CLA	6,4	00054040
	PAC	0,3	00054050
	CLA	7,4	00054060
	PAC	0,5	00054070
A51	TXI	*+1,3,2L -2	00054080
.	TXI	*+1,5,2L -2	00054090
.			00054100
.			00054110
A2			00054120
	AXI	2L,2	00054130
	PXA	0,3	00054140
	PAX	0,6	00054150
	PAX	0,7	00054160
	PXA	0,0	00054170
A3	LDQ	A,6	00054180
	LRS	0	00054190
	TLQ	*+2	00054200
	XCA		00054210
	LDQ	A+1,6	00054220
	LRS	0	00054230
	TLQ	*+2	00054240
	XCA		00054250
	TNX	A4,2,2	00054260
	TXI	A3,6,2N	00054270
A4	TZE	E1	00054280
	STG	T	00054290
	CLA	=1.0	00054300
	FDP	T	00054310
	STQ	T	00054320
	AXI	2L,2	00054330
A5	FMP	A,7	00054340
	STG	A,7	00054350

## NORMALIZATION OF ROWS



# COMPLEX SIMULTANEOUS EQUATION SUBROUTINE

A6	LDQ	T	00054360
	FMP	A+1,7	00054370
	STG	A+1,7	00054380
	LDQ	T	00054390
	TNX	A6,2,2	00054400
	TXI	A5,7,2N	00054410
	PXA	0,5	00054420
	PAX	0,6	00054430
	AXT	LB,4	00054440
1A6	FMP	B,6	00054450
	STG	B,6	00054460
	LDQ	T	00054470
	FMP	B+1,6	00054480
	STG	B+1,6	00054490
	LDQ	T	00054500
2A6	TNX	2A6,4,1	00054510
	TXI	1A6,6,2N	00054520
	TNX	A7,1,2	00054530
	TXI	B+1,3,2	00054540
	TXI	A2,5,2	00054550
A7	TXH	B7,2,2L -2	00054560
.			00054570
.			00054580
.			00054590
	PXA	0,2	00054600
	PAX	0,1	00054610
	PXA	0,3	00054620
	PAX	0,6	00054630
	PXA	0,0	00054640
A8	LDQ	A,6	00054650
	LRS	0	00054660
	TLQ	B+3	00054670
	XCA		00054680
	SXA	A10,1	00054690
	LDQ	A+1,6	00054700
	LRS	0	00054710
	TLQ	B+3	00054720

SEARCH FOR MAXIMUM-PIVOT IN COLUMN

# COMPLEX SIMULTANEOUS EQUATION SUBROUTINE

	XCA			00054730
	SXA	A10,1		00054740
	TXI	+1,1,2		00054750
	TXH	+2,1,2L		00054760
A9	TXI	A8,6,-2		00054770
	LDQ	TOL		00054780
	TLQ	+2		00054790
	TRA	E3		00054800
A10	AXT	+1		00054810
	SXD	+1,2		00054820
	TNX	A17,1,**		00054830
				00054840
				00054850
				00054860
				00054870
				00054880
				00054890
				00054900
				00054910
				00054920
				00054930
				00054940
				00054950
				00054960
				00054970
				00054980
				00054990
				00055000
				00055010
				00055020
A12	TXI	+1,4,2		00055030
	TXH	A13,4,2L		00055040
	TXI	+1,6,2N		00055050
	TXI	A11,7,2N		00055060
A13	PXA	0,5		00055070
	PAX	0,6		00055080
	PAX	0,7		00055090
A14	AXT	LB,4		

	ROW INTERCHANGE
	PXA 0,3
	PAX 0,6
	PAX 0,7
	SCO +1,1
	TXI +1,7,**
	PXA 0,2
	PAX 0,4
	CLA A,6
	LUQ A,7
	STG A,7
	STQ A,6
	CLA A+1,6
	LDQ A+1,7
	STG A+1,7
	STQ A+1,6
	TXI +1,4,2
	TXH A13,4,2L
	TXI +1,6,2N
	TXI A11,7,2N
	PXA 0,5
	PAX 0,6
	PAX 0,7
	AXT LB,4

# COMPLEX SIMULTANEOUS EQUATION SUBROUTINE

A15	SCD	•+1,1	00055100
	TXI	•+1,6,••	00055110
	CLA	8,7	00055120
	LQG	8,6	00055130
	STG	R,6	00055140
	STQ	8,7	00055150
	CLA	8+1,7	00055160
	LQG	8+1,6	00055170
	STG	8+1,6	00055180
	STQ	8+1,7	00055190
	TNX	A17,4,1	00055200
A16	TXI	•+1,6,2N	00055210
	TXI	A15,7,2N	00055220
•			00055230
•			00055240
•			00055250
A17	N2T	A,3	00055260
	TRA	B1	00055270
	LQG	A,3	00055280
	FMP	A,3	00055290
	STG	T	00055300
	LQG	A+1,3	00055310
	FMP	A+1,3	00055320
	FAD	T	00055330
	STG	T	00055340
	CLA	A,3	00055350
	FOP	T	00055360
	STQ	AM	00055370
	CLA	A+1,3	00055380
	FOP	T	00055390
	STQ	AN	00055400
A18	TXH	A21,2,2L -2	00055410
	PXA	O,2	00055420
	PAX	O,4	00055430
	PXA	O,3	00055440
A20	PAX	O,6	00055450
	TXI	•+1,6,2N	00055460

DIVISION OF ROW BY PIVOT

# COMPLEX SIMULTANEOUS EQUATION SUBROUTINE

LDQ	A,6	00055470
FMP	AN	00055480
STO	T	00055490
LDQ	A+1,6	00055500
FMP	AM	00055510
FSB	T	00055520
LDQ	A+1,6	00055530
STO	A+1,6	00055540
FMP	AN	00055550
STO	T	00055560
LDQ	A,6	00055570
FMP	AM	00055580
FAD	T	00055590
STO	A,6	00055600
TXI	+1,4,2	00055610
TXL	A20,4,2L -2	00055620
PXA	0,5	00055630
PAX	0,6	00055640
AXI	LB,4	00055650
LDQ	B,6	00055660
FMP	AN	00055670
STO	T	00055680
LDQ	B+1,6	00055690
FMP	AM	00055700
FSB	T	00055710
LDQ	B+1,6	00055720
STO	B+1,6	00055730
FMP	AN	00055740
STO	T	00055750
LDQ	B,6	00055760
FMP	AM	00055770
FAD	T	00055780
STO	B,6	00055790
TXI	+2,4,1	00055800
TXL	1A21,6,2N	00055810
2A21		00055820
		00055830

ROW REDUCTION

# COMPLEX SIMULTANEOUS EQUATION SUBROUTINE

	PXA	0,2	00055840
	PAX	0,1	00055850
	PAX	0,4	00055860
	TNX	A31,1,2	00055870
	PXA	0,3	00055880
	PAX	0,6	00055890
	PAX	0,7	00055900
	STA	A29	00055910
	SXA	A26+1,5	00055920
	SXA	3A26,5	00055930
	TXI	+1,6,2	00055940
A22	TXI	+1,7,2N -2	00055950
	TXI	+1,3,2N -2	00055960
	SXA	A28,7	00055970
A23	TXH	A26,2,2L -2	00055980
	SXA	A27,3	00055990
A24	LQ	A,6	00056000
	FMP	A,7	00056010
	STG	T	00056020
	LQ	A+1,6	00056030
	FMP	A+1,7	00056040
	FSB	T	00056050
	FAC	A,3	00056060
	STG	A,3	00056070
	LQ	A,6	00056080
	FMP	A+1,7	00056090
	STG	T	00056100
	LUQ	A+1,6	00056110
	FMP	A,7	00056120
	FAD	T	00056130
	CHS		00056140
	FAD	A+1,3	00056150
	STG	A+1,3	00056160
	TXI	+1,4,2	00056170
A25	TXH	A27,4,2L -2	00056180
	TXI	+1,3,2N	00056190
			00056200

# COMPLEX SIMULTANEOUS EQUATION SUBROUTINE

A27	TXI	A24,7,2N	00056210
A26	AXT	0,3	00056220
	AXT	L8,4	00056230
	AXT	0,7	00056240
	TXI	0+1,7,2	00056250
	SXA	0-2,7	00056260
1A26	LDQ	A,6	00056270
	FMP	B,5	00056280
	STO	T	00056290
	LDQ	A+1,6	00056300
	FMP	B+1,5	00056310
	F33	T	00056320
	FAD	B,7	00056330
	STO	B,7	00056340
	LDQ	A,6	00056350
	FMP	B+1,5	00056360
	STO	T	00056370
	LDQ	A+1,6	00056380
	FMP	B,5	00056390
	FAD	T	00056400
	CHS		00056410
	FAD	B+1,7	00056420
	STO	B+1,7	00056430
2A26	TNX	3A26,4,1	00056440
	TXI	0+1,5,2N	00056450
3A26	TXI	1A26,7,2N	00056460
	AXT	0,5	00056470
	TNX	A29,1,2	00056480
A28	AXT	0,7	00056490
	PXA	0,2	00056500
	PAX	0,4	00056510
	TXI	0+1,3,2	00056520
	TXI	A23,6,2	00056530
A29	AXT	0,3	00056540
A31	TXH	A43,2,2L -2	00056550
	PXA	0,2	00056560
	PAX	0,1	00056570

# COMPLEX SIMULTANEOUS EQUATION SUBROUTINE

	PAX	0,4	00056580
	PXA	0,3	00056590
	PAX	0,6	00056600
	PAX	0,7	00056610
A32	TXI	*+1,3,2N +2	00056620
	TXI	*+1,6,-2	00056630
A33	TXI	*+1,7,2N	00056640
	SXA	A37+1,5	00056650
	SXA	A41,5	00056660
	SXA	A40,3	00056670
	SXA	A39,7	00056680
A34	SXA	A38,3	00056690
A35	LDQ	A,6	00056700
	FMP	A,7	00056710
	STG	T	00056720
	LDQ	A+1,6	00056730
	FMP	A+1,7	00056740
	FSB	T	00056750
	FAD	A,3	00056760
	STG	A,3	00056770
	LDQ	A,6	00056780
	FMP	A+1,7	00056790
	STG	T	00056800
	LDQ	A+1,6	00056810
	FMP	A,7	00056820
	FAD	T	00056830
	CHS		00056840
	FAD	A+1,3	00056850
	STG	A+1,3	00056860
	TXI	*+1,4,2	00056870
A36	TXM	A37,4,2L -2	00056880
	TXI	*+1,3,2N	00056890
	TXI	A35,7,2N	00056900
A37	AXT	LB,4	00056910
	AXT	*+7	00056920
	TXI	*+1,7,-2	00056930
	SXA	*-2,7	00056940

COMPLEX SIMULTANEOUS EQUATION SUBROUTINE

1A37	LDQ	A,6	00056950
	FMP	B,5	00056960
	STG	T	00056970
	LDQ	A+1,6	00056980
	FMP	B+1,5	00056990
	FSB	T	00057000
	FAD	B,7	00057010
	STG	B,7	00057020
	LDQ	A,6	00057030
	FMP	B+1,5	00057040
	STG	T	00057050
	LDQ	A+1,6	00057060
	FMP	B,5	00057070
	FAD	T	00057080
	CHS		00057090
	FAD	B+1,7	00057100
	STG	B+1,7	00057110
	TNX	A41,4,1	00057120
2A37	TXI	+1,5,2N	00057130
	TXI	1A37,7,2†	00057140
A41	AXT	+,5	00057150
	TXI	+1,1,2	00057160
	TXH	A40,1,2L -2	00057170
A36	AXT	+,3	00057180
A39	AXT	+,7	00057190
	PXA	0,2	00057200
	PAX	0,4	00057210
	TXI	+1,3,-2	00057220
	TXI	A34,6,-2	00057230
A40	AXT	+,3	00057240
	TXI	+1,5,-2	00057250
	TXI	A7,2,2	00057260
A43	CLA	=1	00057270
	TRA	MSIMEC+1	00057280
B1	STZ	AM	00057290
	CLA	=1.0	00057300
	FDP	A+1,3	00057310





#### APPENDIX IV. SAMPLE DATA SHEETS

The following pages are sample data sheets for a computer run on three modes at three frequencies. The potential will not be computed for the first mode. The generalized forces four will be  $L_{21}$ ,  $L_{22}$ ,  $L_{23}$ ,  $L_{31}$ ,  $L_{32}$ ,  $L_{33}$ .

Of the first fourteen cards, the cards numbered 6, 9, 10, 11, 12, 13, 14 do nothing and are included only to indicate how all data is entered. Cards 1 through 14 are complete in this respect, and all later cards are of the same type as one of the first fourteen. The data used in the least squares surface fit for the deflection is entered in locations 101 through 700.

Card number 22 represents 56 cards for the intermediate data points which would have to be included in an actual run.

# FORTRAN FIXED 10 DIGIT DECIMAL DATA

DECK NO.	PROGRAMMER	DATE	PAGE	of	JOB NO.
1	A L P H A				
2	6 0				
3	P P E U D E I L T A				
4	A T 2 0 C P S				
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7	A L P H A				
8	P L U N G E				
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12	2 3				
13	2 0 . 0				
14	2 5 . 0				
15	1 2 8 7 5 . 0				
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# FORTRAN FIXED 10 DIGIT DECIMAL DATA

DECK NO. _____		PROGRAMMER _____		DATE _____		PAGE _____ of _____		JOB NO. _____	
NUMBER		IDENTIFICATION		DESCRIPTION DO NOT KEY PUNCH					
1	3 0	[Shaded Box]		normalization of points on leading edge					
2	0 . 0			76					
3	3 5 . 6 0			77					
4	9 . 0			78					
5		0 5		used if $MB \geq 2$					
6		[Shaded Box]		used if $MB = 3$					
7	3 5			79					
8				80					
9				81					
10		0 6		used for value suppression calculation of potential for this mode					
11	3 9	[Shaded Box]							
12	1								
13									
14									
15		0 7							
16	1 6	[Shaded Box]		beginning of list of coefficients of deflection series					
17	1 . 0			82					
18				83					
19				84					
20		0 8		85					

# FORTRAN FIXED 10 DIGIT DECIMAL DATA

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# FORTRAN FIXED 10 DIGIT DECIMAL DATA

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last card for second mode  
 this entry sample maximum value of 4.0  
 the polynomial is 0.042, normalized for maximum value 1.0

beginning of data for third mode

# FORTRAN FIXED IO DIGIT DECIMAL DATA

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NUMBER		IDENTIFICATION	DESCRIPTION DO NOT KEY PUNCH			
1	1 0 1	2 1	data for first point			
2	2 3 . 6					
3	0 . 5					
4	0 . 0 2 3 5					
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**FORM 1041-13 REV. 2001**



# FORTRAN FIXED 10 DIGIT DECIMAL DATA

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data for a third frequency